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LOKI WIND CORRECTION COMPUTER

by

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and

WIND STUDIES FOR LOKI

by

Bernard Helfand

August 1, 1953

**Annual Summary for June 24, 1952 to June 24, 1953
and Report for Fourth Quarter
April 1 through June 30, 1953**

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Annual Summary

The work performed during the first year of the LOKI Wind Correction Computer development program has fallen into four phases, described in detail in the North American Instruments, Inc. Quarterly Report series.

In the first phase, the effect of an arbitrary wind distribution on the LOKI trajectory during boost was determined by an impulse response method, and a function was derived relating burnout deviation of flight path with the distributed wind along the boost trajectory. Complete details of the derivation are included in the First Quarterly Report, Ref. 5. To verify the theoretical influence function, arrangements have been made for an experimental comparison at White Sands Proving Ground. The Signal Corps is erecting ten poles instrumented with anemometers at the launching site of the Small Missiles Range, for this purpose. This system is expected to be in operation within the next few months.

In the second phase, the wind influence function was then applied to arbitrary and observed wind distributions with a view to establishing the basis for a practical field wind measuring system to transmit information to a computer which could perform an aiming correction. The studies indicated the feasibility of a correction method

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based on a linear relationship between burnout deviation and wind, suitably measured, at a limited number of fixed stations. This work is described in the Second Quarterly Report. The conclusions are supported by further studies included in this report.

To supplement immediately the wind data obtained from the literature, which extend up to 250 feet altitude (50% of the burnout deviation effect), with measurements extending to 1,000 feet of altitude (94% of the burnout deviation effect), a balloon-borne hot wire anemometer system was constructed and put into field operation, data being gathered in the Mojave Desert at El Mirage Airport, and at White Sands Proving Ground at the Small Missiles Range. The equipment is described in the Third Quarterly Report. A more complete program involving space and time variations of wind within a volume significant in the general rocket aiming problem has been set up under the sponsorship of the Signal Corps. The system involves the erection of a number of high towers and the operation of a complex data recording system. It is expected that data will become available from this project within the next year.

At the outset of the computer development program the possible need was recognized for an anemometer capable of fast response, to provide orthogonal wind components in terms of electrical outputs. No such instrument was available and a design was undertaken based on the drag force exerted by the wind stream on a transverse cylinder.

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The development of this 2-component anemometer constituted the third phase and extended through most of the year and resulted in a prototype unit suitable for field operation. Details on the design and wind tunnel testing of the anemometer are contained in the Third Quarterly Report.

The final phase of the past year's work, described in this report, has been the design of a complete wind correction computer system based on the influence function and the results of the wind studies.

Acknowledgment

We wish to acknowledge helpful discussions, held during the preparation of this report, with members of the Bell Telephone Laboratories on the subjects of Computer Geometry and Components.

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PART I

WIND CORRECTION COMPUTER

Introduction

The purpose of the computer which is the subject of this report is to correct the aim of a rocket launcher (in both azimuth and elevation) for the wind forces experienced by the rocket during its burning period. As a result of a statistical study of winds within a few hundred feet of the ground, it is believed that the effect of these so-called surface winds upon the rocket trajectory can be expressed by a linear relation involving the two horizontal wind components at or near the rocket launcher. The details of the wind structure study are presented in this report and in an earlier one, Ref. 1, and the anemometer used in both the wind study and in the wind correction computer has been described earlier, Ref. 2.

The angular deviation at the end of burning due to wind normal to the trajectory, amounts to approximately one mil (thousandth radian) per mile-per-hour; it is not expected that rocket launching will be attempted in winds above 40 miles per hour. Current experimental rocket dispersions amount to 8 mils (linear standard deviation), and a wind correction standard deviation of about 2 mils thus appears acceptable.

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Operating Principle of Computer

The wind correction computer corrects the point of aim of the rocket launcher for deviations from its intended trajectory due to wind. The corrections are made in both azimuth and elevation; both corrections are inserted as rotations of the shafts of differential selsyns interposed in the transmission lines between the main fire control director (M33) and the launcher. The wind correction computer is physically separate from the main fire control director and may be remote from it.

The wind data received from the anemometer is in the form of two amplitude-modulated 3000 cps signals; the amplitudes are proportional to $V^2 \cos \Theta$ and $V^2 \sin \Theta$, the northerly and westerly components of the square of the wind velocity; the wind is assumed to lie in the horizontal plane. Here V is the wind velocity and Θ is the wind vector heading angle measured counter-clockwise relative to true north. Because the anemometers do not deliver actual velocity components, the computer must convert $V^2 \cos \Theta$ and $V^2 \sin \Theta$ into $V \cos \Theta$ and $V \sin \Theta$, respectively. Secondly, if the anemometer is fixed in orientation, (i. e., not slaved in azimuth to the launcher), it is necessary to carry out the vector rotations which resolve the northerly and westerly wind components into down-range and cross-range winds; this is achieved in the computer by sine and cosine potentiometers and operational adders. If A is the azimuth

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angle of the launcher, the desired quantities are:

$$V \cos(\Theta - A) = (V \cos \Theta) \cos A + (V \sin \Theta) \sin A \quad (1)$$

$$V \sin(\Theta - A) = -(V \cos \Theta) \sin A + (V \sin \Theta) \cos A \quad (2)$$

We may replace $(\Theta - A)$ by ϕ , the wind heading relative to the launcher azimuth.

Since only those wind components normal to the trajectory are effective in producing deflections of the trajectory, the angular rotation due to wind may be resolved into two angular rotations normal to the trajectory. The horizontal and vertical angular deviations W_h and W_v , respectively, may be obtained by the geometric relations, illustrated in Fig. 1. Because of the smallness of the angular deviation, the deviation resolves in the same manner as the wind force which causes it. The incremental angles are:

$$W_h = KV \sin \phi \quad (3)$$

$$W_v = KV \cos \phi \sin E \quad (4)$$

where E is the corrected angle of launcher elevation and K is the influence coefficient relating wind (in mph) to angular deviation (in mils).

It should be remembered that the deviations are angular and the incremental vectors, W_n , W_h , and W_v are chords of a unit sphere. Fig. 2 illustrates the spherical trigonometric relations utilized in deriving the wind corrections in azimuth and elevation.

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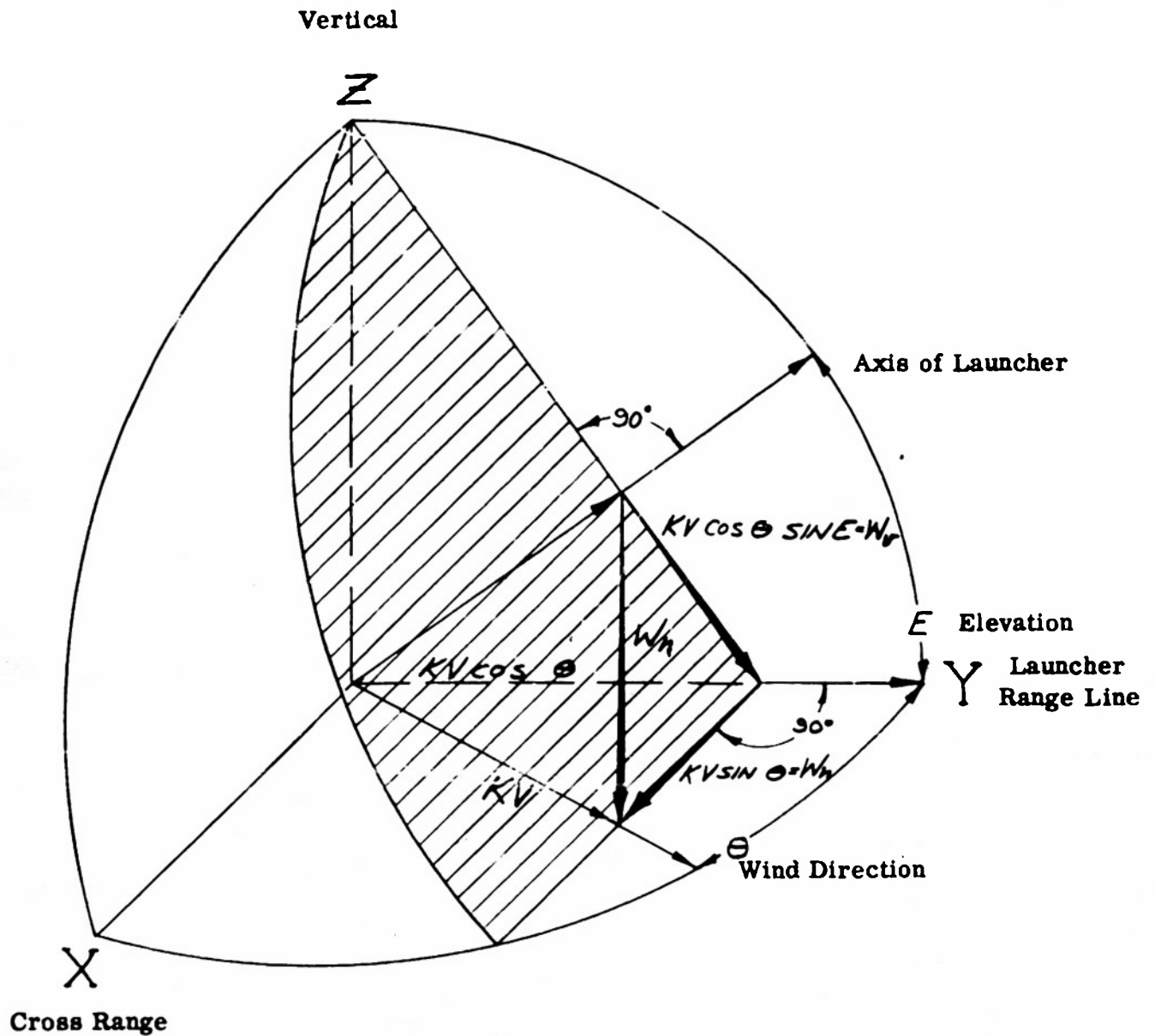


FIG. 1 RESOLUTION OF WIND VECTOR OR ANGULAR DEVIATION

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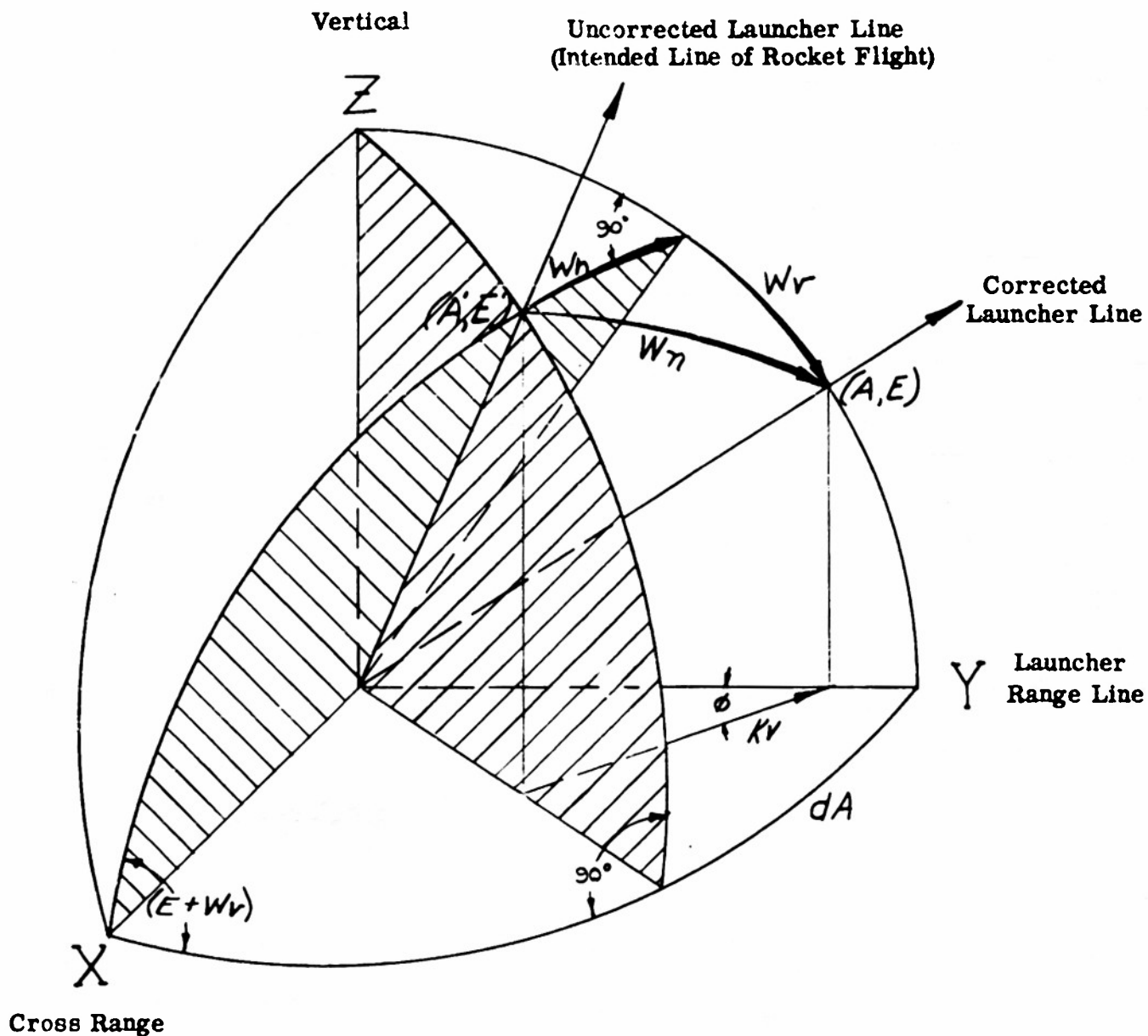


FIG. 2 RESOLUTION OF ANGULAR DEVIATION DUE TO WIND INTO AZIMUTH (A) AND ELEVATION (E) CORRECTIONS. PRIMED QUANTITIES ARE UNCORRECTED OR "AIM-POINT" VALUES.

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By the law of sines, one may write

$$\sin dA = \sin W_h / \cos E' \quad (5)$$

The quantity $\sin W_h$ may be replaced by W_h to within one part in 4,000; $\sin dA$ cannot be replaced by dA because dA may be as great as 560 mils; thus the equation for dA may be written:

$$\sin dA = W_h \sec E' \quad (6)$$

The correction in elevation, dE , may be written by a second application of the law of sines as:

$$\sin (E + W_v) = \sin E' / \cos W_h \quad (7)$$

We shall want the expression in a form not involving E , so we substitute

$$E = E' + dE \quad (8)$$

with the result:

$$dE = -W_v + \frac{W_h^2}{2} \tan E' \quad (9)$$

In writing Equation (9), we have made use of the approximations:

$$\cos (dE + W_v) \approx 1 \quad 0.3\% \quad (10)$$

$$\sin (dE + W_v) \approx (dE + W_v) \quad 0.1\% \quad (11)$$

$$\cos W_h \approx 1 - \frac{W_h^2}{2} \quad 10^{-5}\% \quad (12)$$

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Equations (3), (4), (6), and (9) provide the mathematical basis for the computation of wind corrections. Equation (4) involves the angle E while all other relations involve E'.

If we denote

$$W'_v = KV \cos \phi \sin E' \quad (13)$$

$$\text{then } (1 + W'_v \cot E') dE = -W'_v + \frac{W_h^2}{2} \tan E' \quad (14)$$

Because of the small magnitude of W'_v , neglect of the term $W'_v \cot E'$ causes an error of less than 4% in dE. Hence, to within 1.6 mils, Equation (9) may be replaced by

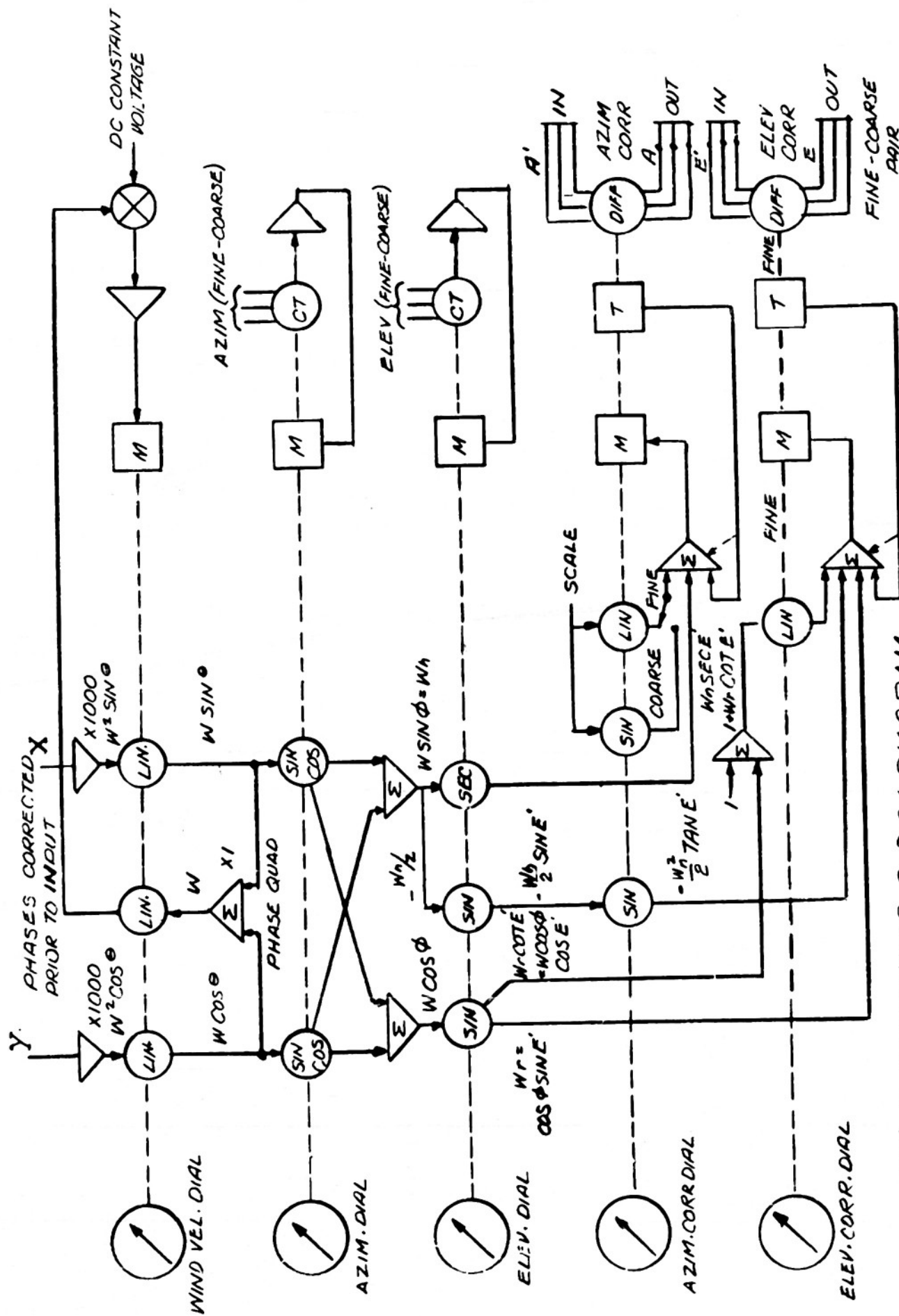
$$dE = -W'_v + \frac{W_h^2}{2} \tan E' \quad (15)$$

the corrected elevation, E, does not appear in any of the computer relations and no servo shaft positioning device need be provided for it. If desirable, the computer can be arranged to avoid this slight error. Fig. 3 indicates a solution utilizing Equation (14).

Computer Block Diagram

Fig. 3 presents a block diagram of the wind correction computer. The five horizontal dashed lines represent mechanical shafts, positioned by servomotors on the right; several linear or functional potentiometers are placed on each shaft; indicating dials (and selsyn or potentiometer transmitters, if required) are on the left side of the diagram.

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The manner in which the computer executes the mathematical relations may be seen fairly readily. The two anemometer component signals $W^2 \sin \Theta$ and $W^2 \cos \Theta$, are amplified to suitable level and introduced at the top of Fig. 3. The three linear potentiometers and the quadrature superposition unit are used for extracting the square root of V^2 and dividing each of the input signals to yield $W \cos \Theta$ and $W \sin \Theta$. These signals are combined by means of the azimuth sine-cosine potentiometers and subsequent adding amplifiers to produce the components of the correction relative to the instantaneous launcher heading, $W \cos \phi$ and $W \sin \phi$.

Several functional potentiometers are employed to modulate the $V \cos \phi$ and $V \sin \phi$ signals by $\sin E'$, $\cos E'$, and $\sec E'$; the outputs from these functional potentiometers are combined in several ways by adding amplifiers and the resulting signals are balanced by output servos against the signals derived from other potentiometers mounted on the output shafts.

Square-Root Techniques

The potentiometer cascade utilized in the computer schematic of Fig. 3 is rather simple in principle, but may lead to significant dynamic error in operation due to the slowness of the servo-balancing as compared to wind fluctuations. An analysis of the dynamic error due to this cause has not yet been carried out; however, it is apparent that the error

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may amount to several percent unless the input signals are damped to correspond to the servo capabilities.

Damping may be achieved by means of an orifice in the anemometer (up to about 15 seconds time constant) or by filter networks placed in d. c. sections of the computer; experimentally, the latter method is more appealing because of the possibility of switching damping values. However, it is convenient to avoid the troubles of d. c. amplification in the computer by using the 3000 cps signals throughout; this favors mechanical damping. To lessen the importance of this controversy, it is contemplated to make the experimental computer convertible insofar as possible from a. c. to d. c. signal operation; this is the reason reactive resolvers are not used. If d. c. operation is employed, a. c. signals are none the less required through the quadrature circuits, and the rectification must take place after the square root has been extracted.

To eliminate the possible source of dynamic error in square rooting, two non-servo-square-root circuits have been investigated. One type of square-rooter utilizes the non-linear properties of Thyrite, while a second type employs suppressor controlled pentodes in a feedback circuit to extract the square root; simplified breadboard computers have been built utilizing both approaches. The two methods of square rooting are treated further in the following section.

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Thyrte Type Computer

The computer utilizing Thyrte as a square root extracting device combines the signals $V^2 \cos \theta$ and $V^2 \sin \theta$ in phase quadrature to produce a signal of amplitude V^2 . The Thyrte square root extractor delivers d. c. voltages proportional to $+|V|$ and $-|V|$, which are applied to the two potentiometers. The output angular positions representing the azimuth and elevation corrections are positioned by servo-mechanisms effecting the following relations:

$$|V| \sin \Delta A = KV^2 \sin \phi \sec E' \quad (16)$$

$$|V| \sin \Delta E = KV^2 \cos \phi \sin E' \quad (17)$$

The computer requires the solution for which $|V| \neq 0$; for this reason $|V|$ is constrained to be significantly greater than zero at all times, regardless of wind. The use of a variable voltage $|V|$ across the servo potentiometers leads to a system in which the servo loop gain is variable; as is well known, this situation is usually conducive to instability under some of the desired operating conditions. This situation might be tolerated in this mode of operation because of the convenience and economy of effecting both the square root extraction and the output positioning with the same servo.

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Thyrrite Characteristics

The non-linear characteristics of Thyrrite, a material manufactured by the General Electric Company for use in lightning arrestors, are demonstrated by the plot of Fig. 4; the graph shows the voltage across a small Thyrrite element as a function of the current passing through it. The slight departure from the desired square-root dependence is almost entirely corrected by the use of a trimming resistor, in this case about 300 ohms.

A very simple Thyrrite driver circuit was tried, which requires only moderate voltages; the circuit is shown in Fig. 5 and its performance is represented by Fig. 7. The 5% accuracy achieved in this trial circuit can undoubtedly be improved to correspond to the 2-3% result of the Thyrrite alone.

Suppressor Type Computer

To overcome the necessity for d. c. servo amplification and to avoid the hazards of variable servo-loop gain, some tests have been made of square root extraction by means of successive suppressor controlled variable gain amplification stages. The Suppressor Type Computer utilizes this type of circuit for square root extraction. The computer shown in Fig. 3 is different from the Suppressor Type only in its use of two successive potentiometer stages for root extraction rather than suppressor controlled amplification

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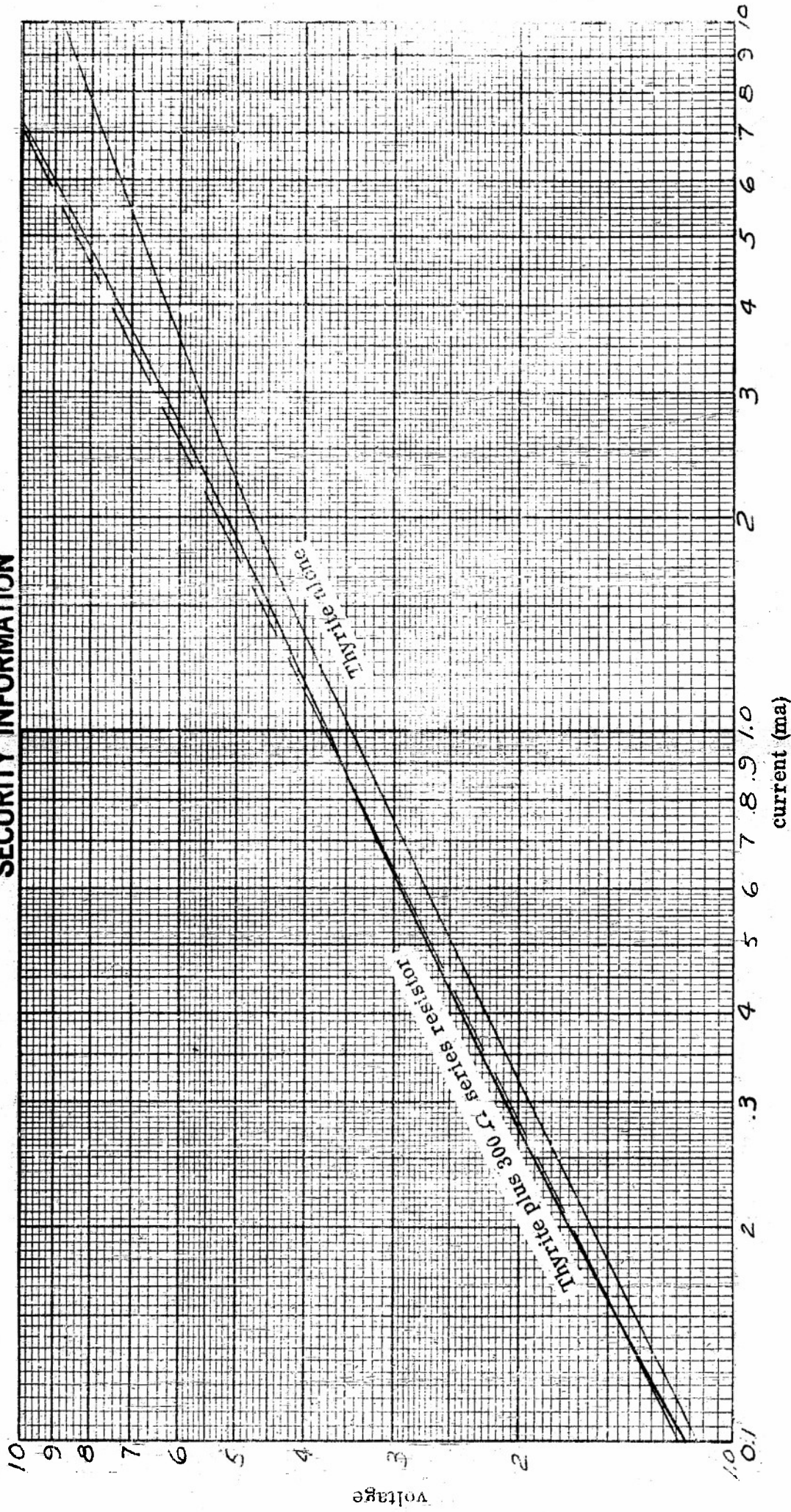


FIG. 4 THYRITE CHARACTERISTICS

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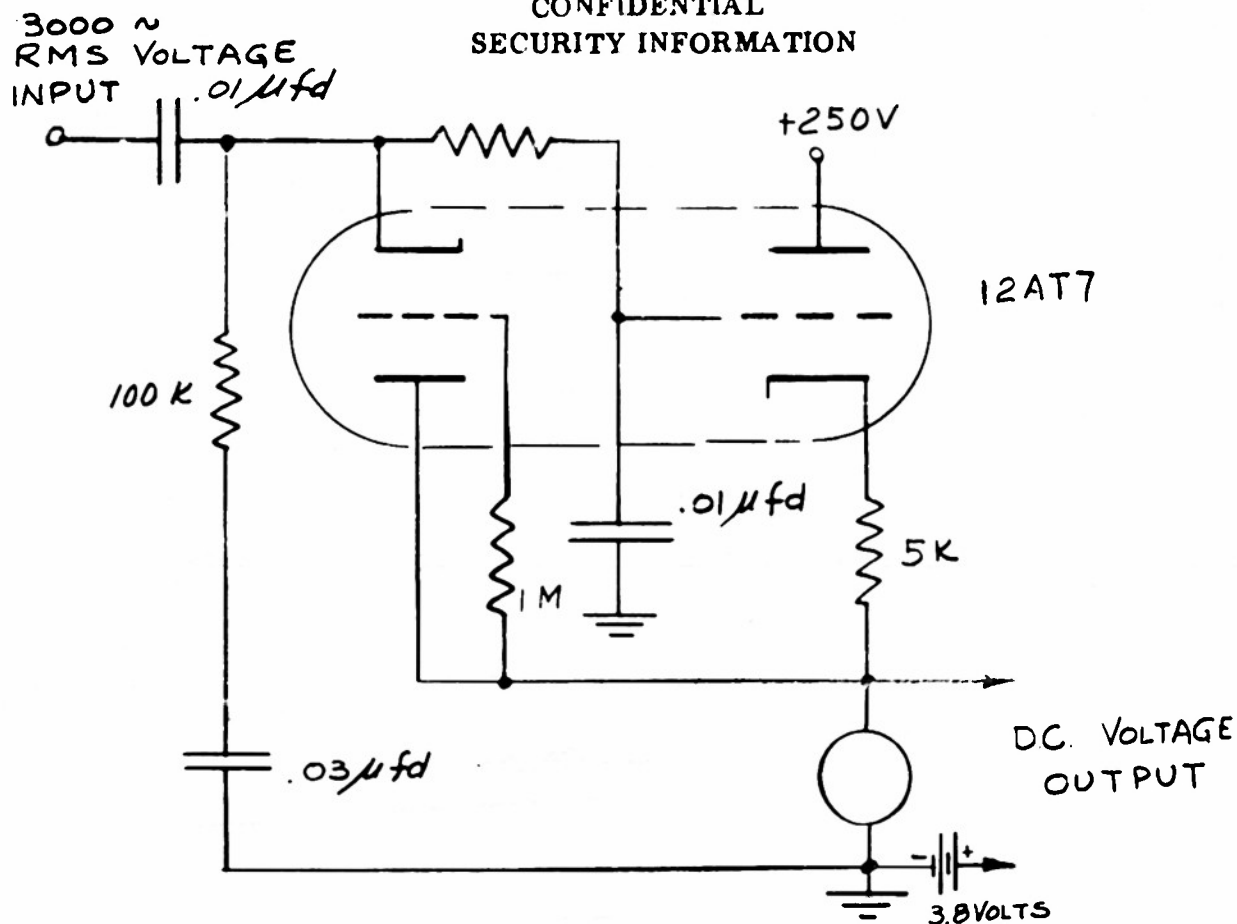
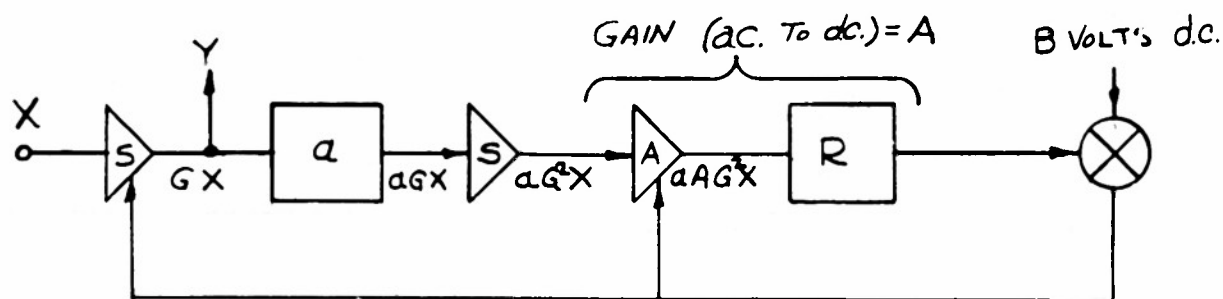


FIG. 5 THYRITE DRIVER CIRCUIT



$$\theta = B - aAX(K + \delta\theta)^2$$

G = GAIN OF THE SUPPRESSORS

$$G = (K + \delta\theta)$$

$$\theta = B - aAX(K + \delta\theta)^2$$

$$\therefore K + \delta\theta = K + \delta B - aA \delta X (K + \delta\theta)^2$$

$$\text{OR } aA \delta X G^2 + G - (K + \delta B) = 0$$

$$\therefore G = \frac{-1 + \sqrt{1 + 4aA\delta(K + \delta B)X}}{2aA\delta X}$$

FIG. 6 PRINCIPLE OF SQUARE ROOT EXTRACTION WITH SUPPRESSOR AMPLIFIERS

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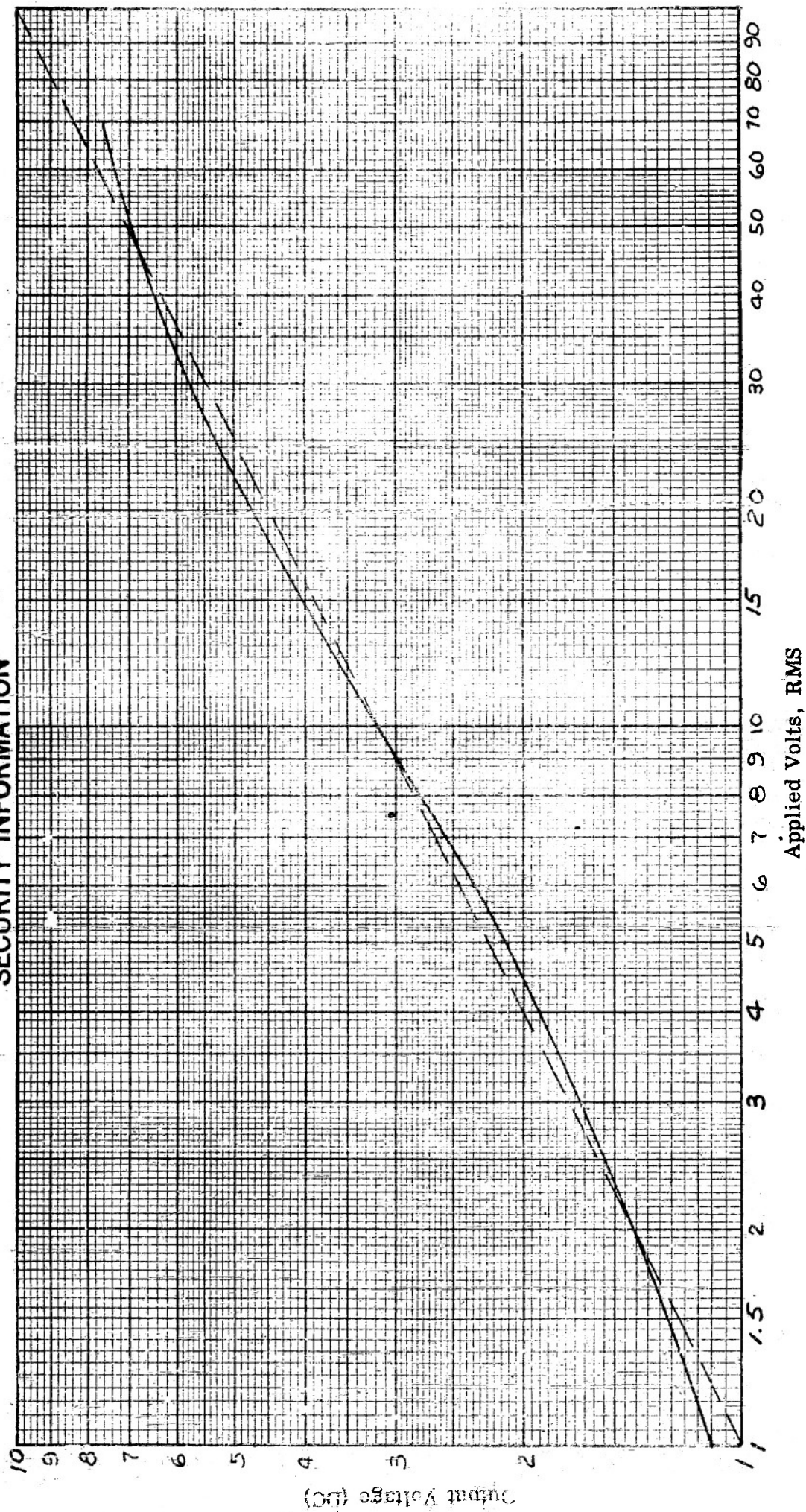


FIG. 7 THYRITE DRIVER PERFORMANCE

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stages. This type of computer, whether it uses suppressor amplifiers or potentiometers for root extraction, has the advantage that no d. c. amplification is required; only simple phase sensitive rectification is required at the output stages of the servo amplifiers.

Suppressor Amplifier Root Extraction

The principle of square root extraction with a feedback circuit incorporating suppressor amplifiers is outlined in Fig. 6. The amplifiers labelled "S" are stages employing, for example, 6AS6 pentodes in which the gain can be controlled readily by the suppressor grid bias. Assuming linear characteristics in the range of operation, the gain of each stage is

$$G = \frac{-1 + \sqrt{1 + 4(K + \delta B) \delta aAx}}{2aA\delta x} \quad (18)$$

where K is the gain of the suppressor amplifier at zero bias, δ is the change of gain per volt of bias, aA is the constant portion of the loop gain (subsequent to the suppressor amplifiers and including the rectification), and X is the input signal (RMS volts). A has the units d. c. volts per RMS volt. B is the d. c. voltage that is compared with the output of the rectifier. The output of the first suppressor stage is, therefore:

$$Y = GX \approx \sqrt{\frac{(K + \delta B)x}{aA}} \left(1 + \frac{1}{8aA \delta (K + \delta B)x} \right) - \frac{1}{2aA \delta} \quad (19)$$

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Typical values of the parameters are:

$$K = 65$$

$$\delta = 6.2 \text{ per volt}$$

$$g_A = 67$$

$$B = 50 \text{ volts d.c.}$$

The calculated behavior of this system shows an error equivalent to less than 0.5 mil; experimental results are not yet available.

The suppressor amplifier square root scheme assumes that the gains of the several suppressor amplifiers vary in unison; to the extent that this is required it constitutes a strong reason against the use of the scheme in an operation computer. However, it can be shown that compensation for differences in tube characteristics can be achieved in large measure by interchanging tubes so that, for example

$$G_2 = G_1 [1 + \alpha(x)], \alpha(x) \geq 0 \quad (20)$$

and adjusting the loop gain so that, approximately

$$G_1 G_2 X = C [1 + \alpha(x)] \quad (21)$$

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PART II

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WIND STUDIES FOR LOKI

Introduction

In an earlier report of this project, Ref. 1, a series of vertical wind profiles was studied to determine the feasibility of a wind correction computer for LOKI, based on a linear relationship between burnout deviation and wind at a selected level. Only a limited number of wind profiles was involved in the comparison. However, significant correlations were found between computed burnout deviation and corresponding single level wind measurements. The effects of altitude and smoothing were examined, comparisons being made with wind level and time averaged velocity as parameters.

The wind data used was that of Sherlock and Stout for winter storms in Michigan, Refs. 3 and 4, and was obtained with a 250-foot tower instrumented with fast response pressure plate anemometers at 25-foot intervals. It was learned from the authors that the published information was only a small part of a large total which had been reduced from the original oscillograph records to punched IBM cards. Several thousand of these cards were available; duplicates were obtained and these were processed by the Los Angeles IBM Service Bureau to extend the initial investigation. The results of the IBM computations which are presented in this report, confirm the initial findings and permit a selection of optimum altitude and time average through the use of a wider range

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of these parameters. Also comparisons were made using combinations of wind values at two altitudes. These appear to offer little improvement over appropriately smoothed single altitude winds at levels above 50 feet.

Discussion

The function relating LOKI deviation at burnout with the distributed wind along the trajectory was derived in Ref. 5 and is shown in Fig. 8 as a plot between wind influence coefficient (deviation of flight path at burnout from direction at launching per unit crosswind per unit distance increment along trajectory) and distance of rocket from launcher. The cumulative effect of a constant crosswind during boost is shown in Fig. 9. The total deviation for unit crosswind is .8 mil, 50% being accumulated in the first 200 feet of trajectory and 95% in the first 1,000 feet.

The results of the preliminary investigation are tabulated in Fig. 10 from Ref. 1. The burnout deviations for 150 cases were calculated by integrating the wind observations weighted with the influence function and apply to a theoretical vertical shot. Since the wind observations extend to 250 feet, the deviation is about 50% of the total that would be obtained if the integration were carried to burnout.

The work discussed in this report extends the number of cases from 150 to 4,000 using the same altitudes, 25, 50, 100, and 150 feet with the addition of the mean of the 25 and

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FIGURE 8
LORI CROSSWIND INFLUENCE FUNCTION

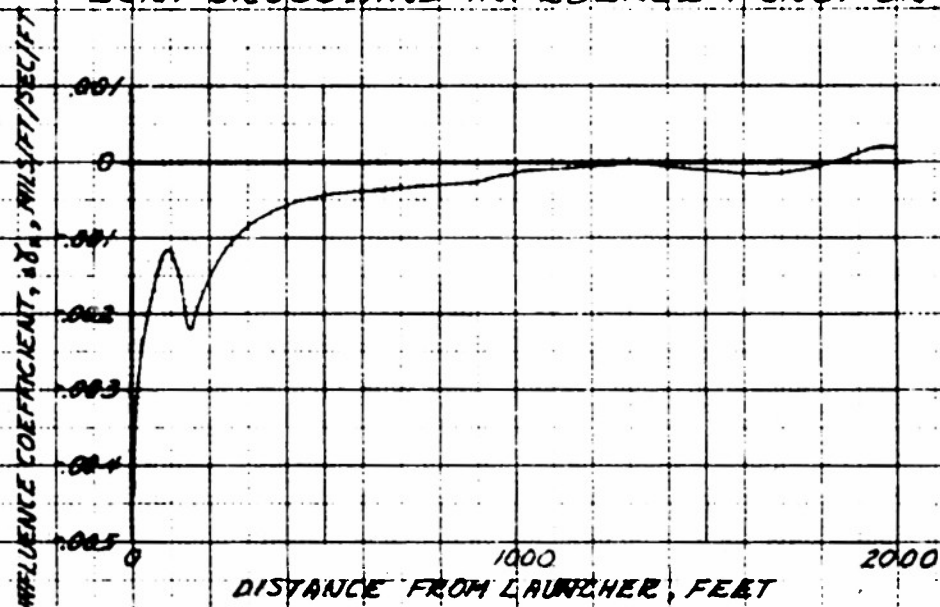
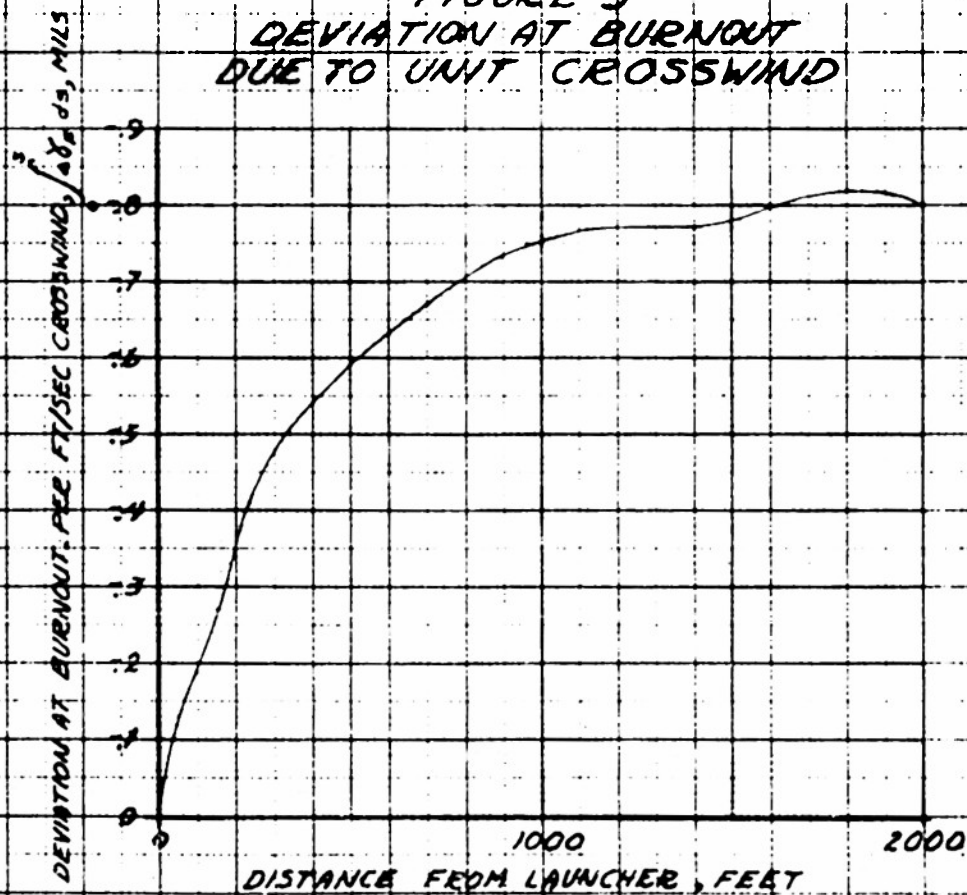


FIGURE 9
DEVIATION AT BURNOUT
DUE TO UNIT CROSSWIND



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Wind Measurement Level	Wind Deviation at Burnout			4-Second Average Wind Deviation at Burnout		
	Regression Line Slope mils/mph	Correlation Coefficient	Deviation Spread for 90% Obs. mils	Regression Line Slope mils/mph	Correlation Coefficient	Deviation Spread for 90% Obs. mils
25	.544	.802	± 4.9	.582	.823	± 4.4
50	.399	.805	± 4.7	.419	.836	± 3.9
100	.488	.901	± 3.3	.491	.896	± 3.2
150	.441	.939	± 2.6	.446	.945	± 2.3

**FIG. 10 SUMMARY OF RESULTS OF PRELIMINARY WIND STUDIES FOR LOKI
(Ref. 1)**

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100-foot winds, and the mean of the 50 and 150-foot winds. Also the 1 and 4-second wind averages in the initial survey are augmented by the 2 and 8-second averages.

The basic wind data as received from the University of Michigan, was in the form of punched IBM cards and apply to the storms of Apr. 28, 1931 and Jan. 19, 1933. All records were originally taken continuously on oscillograph paper, information being transferred from these records onto the cards. The card code identification is shown in Fig. 11.

The first vertical column at the left identifies the run number. Each storm sample was divided into runs of various lengths, the oscillograph being operated intermittently to record during periods of major gust activity. Many runs were found to be broken, due to missing card data, and in order to provide continuous time sequences for analysis, these runs were broken into sub-groups and labelled episodes.

The second column identifies the length of time interval over which the velocity observations were averaged. Averages from 1/4 sec. to 10 sec. were available; however, the averages used in all computations were based on the 1-second cards.

The third column shows the number of seconds after the start of a run, and columns 5 to 16 show the wind speeds in miles per hour at each of the observation stations. The ten stations on the tower were at 25-foot intervals and are identified by numbers 1 to 10 reading from the bottom.

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Code Identification				Velocity in MPH											
				Station Number											
				1	2	3	(4)	5	6	7	8	(9)	10	11	12
Run Number	025	1	Card Number 4 Int. 1/4 Sec.	22.5	25.3	23.9	000	25.1	28.7	27.9	27.8		32.2		
	025	1	Seconds after start of Run	22.9	24.4	23.9	000	25.4	31.1	29.4	30.1		33.5		
	025	1	Time = 1/4 Sec.	etc.	etc.	etc.	000	26.8	29.7	29.0	31.0		33.2		
	025	1									31.9		33.8		
	025	1									32.8		33.5		
	025	1									36.0		33.8		
	025	1									33.3		34.1		
	025	1									32.2		33.8		
	025	1		20.4	24.9	29.3	000	29.4	27.5	28.2	31.6		36.8		
				etc.											

FIG. 11 IBM CARD IDENTIFICATION CODE

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Statistical comparisons were made for each run separately to determine the correlation coefficient, the slope of the burnout deviation vs. wind speed regression line and the 90% confidence limits for the burnout deviation. In calculating the burnout deviation sequences, the wind values at each altitude were assumed to describe the wind extending from the adjacent lower altitude using weighting factors determined from the influence function. The results as tabulated by IBM are presented in the Appendix.

The formulae used in the computations follow:

$$\text{Correlation coefficient, } r = \frac{N \sum \delta A - \sum A \sum \delta}{\{[N \sum A^2 - (\sum A)^2][N \sum \delta^2 - (\sum \delta)^2]\}^{1/2}}$$

$$90\% \text{ Confidence Limits, } Sy_{90} = \frac{1.65 \{[1 - r^2][N \sum \delta^2 - (\sum \delta)^2]\}^{1/2}}{N}$$

$$\text{Regression Line Slope, } b = \frac{N \sum \delta A - \sum \delta \sum A}{N \sum \delta^2 - (\sum \delta)^2}$$

$$\bar{A} = \frac{\sum A}{N}$$

$$\bar{\delta} = \frac{\sum \delta}{N}$$

(Regression Line passes through \bar{A} , $\bar{\delta}$)

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$$\text{Mean Wind } A_{jmk} = \frac{1}{m} \sum_{j=m+1}^j W_{jpk} \quad , \text{ mph}$$

$$\text{Burnout Deviation, } S_{jk} = K \sum_{i=1}^9 \Delta_{ik} W_{ijk} \quad , \text{ miles}$$

where

K is the episode number

j is the time in seconds after the start of a run

m is the index for moving averages, $m = 1, 2, 4, 8$

W is the velocity in miles per hour

Δ_{jk} are the wind weighting factors (see table below)

K is a conversion constant, $K = 88/60$

i is the height of observation given in the following table:

i	Height(feet)	Wind weighting factor Δ_{jk}
1	25	.050
2	50	.050
3	75	.030
4	100	.035
5	125	.030
6	150	.045
7	175	.050
8	200	.040
9	250	.065

p is a specified value or combination of i's as given in the following table:

$$\begin{array}{ll} W_{p_1} = W_{i_1} & W_{p_4} = W_{i_6} \\ W_{p_2} = W_{i_2} & W_{p_5} = \frac{W_{i_1} + W_{i_4}}{2} \\ W_{p_3} = W_{i_4} & W_{p_6} = \frac{W_{i_2} + W_{i_6}}{2} \end{array}$$

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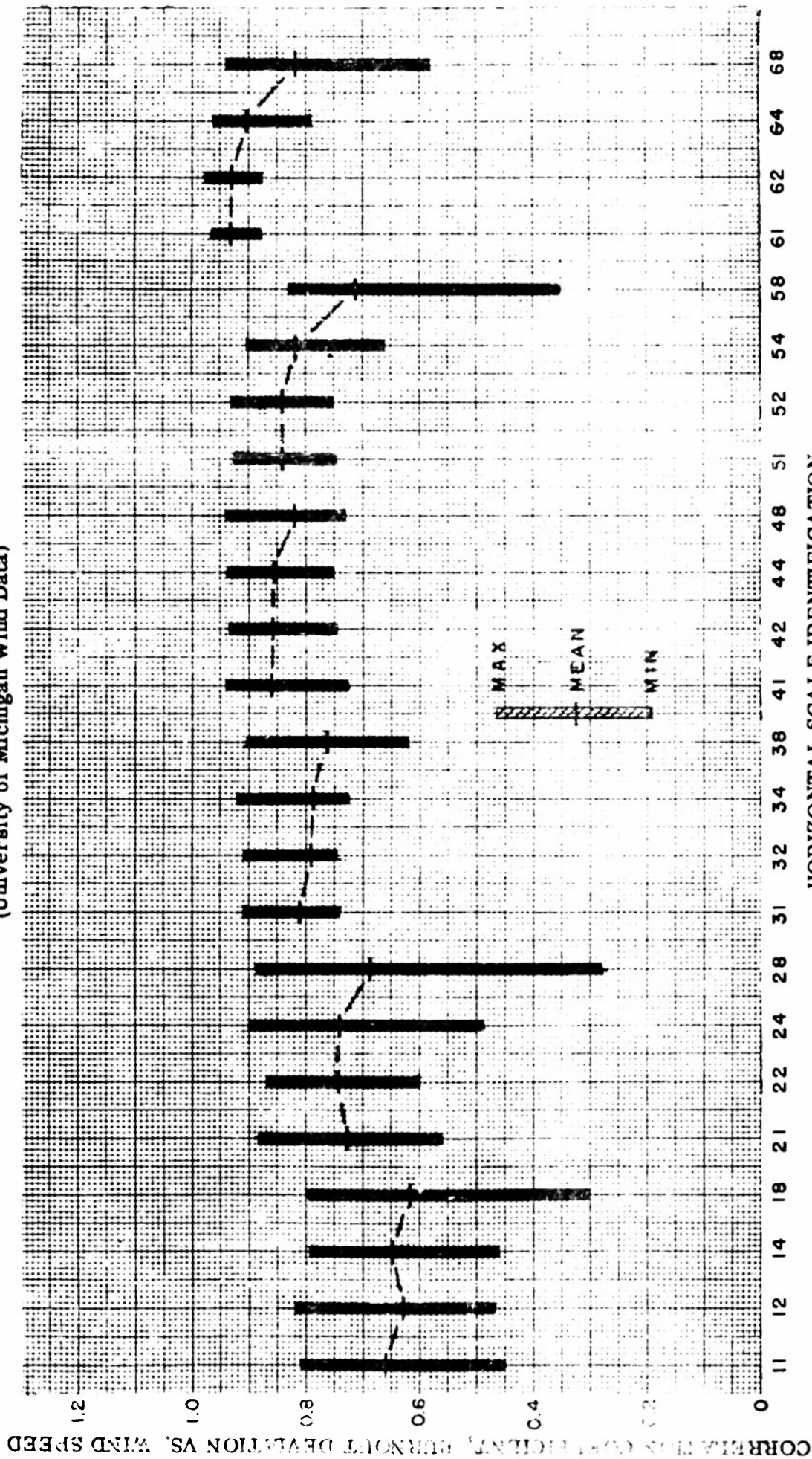
Results are summarized in Figs. 12 through 14, which show the range and average of the three significant parameters plotted as vertical bars and each of the combinations of anemometer height and time average. Runs 17, 19, 34, and 35 have been excluded in certain cases where missing data have resulted in small sample size and resulting range extremes. The unweighted arithmetical mean is shown on each bar. Fig. 12 shows the effect of progressive increase of time average on correlation coefficient of burnout deviation vs. wind speed for progressively increasing anemometer heights from 25 to 150 feet and also for time average 25 to 100-foot winds and 50-150-foot winds. Similar progressive sequences are shown for the regression line slope in Fig. 13 and 90% confidence limits in Fig. 14.

It is evident that the 2-second average offers improved correlation and reduced spread among runs under the 1 and 4-second averages. Further, a pronounced degradation is observed with the 8-second average.

In progressing upward in altitude a marked improvement occurs in going from 50 to 100 feet with subsequent slight improvement at 150 feet. Best results are observed in the 50-150 foot mean wind sequences. However, the improvement over the 100-foot wind would hardly warrant the use of 2 anemometers in a wind correction computing system.

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FIG. 12 EFFECT OF ALTITUDE AND AVERAGE ON CORRELATION COEFFICIENT,
LOKI BURNOUT DEVIATION VS. WIND SPEED
(University of Michigan Wind Data)



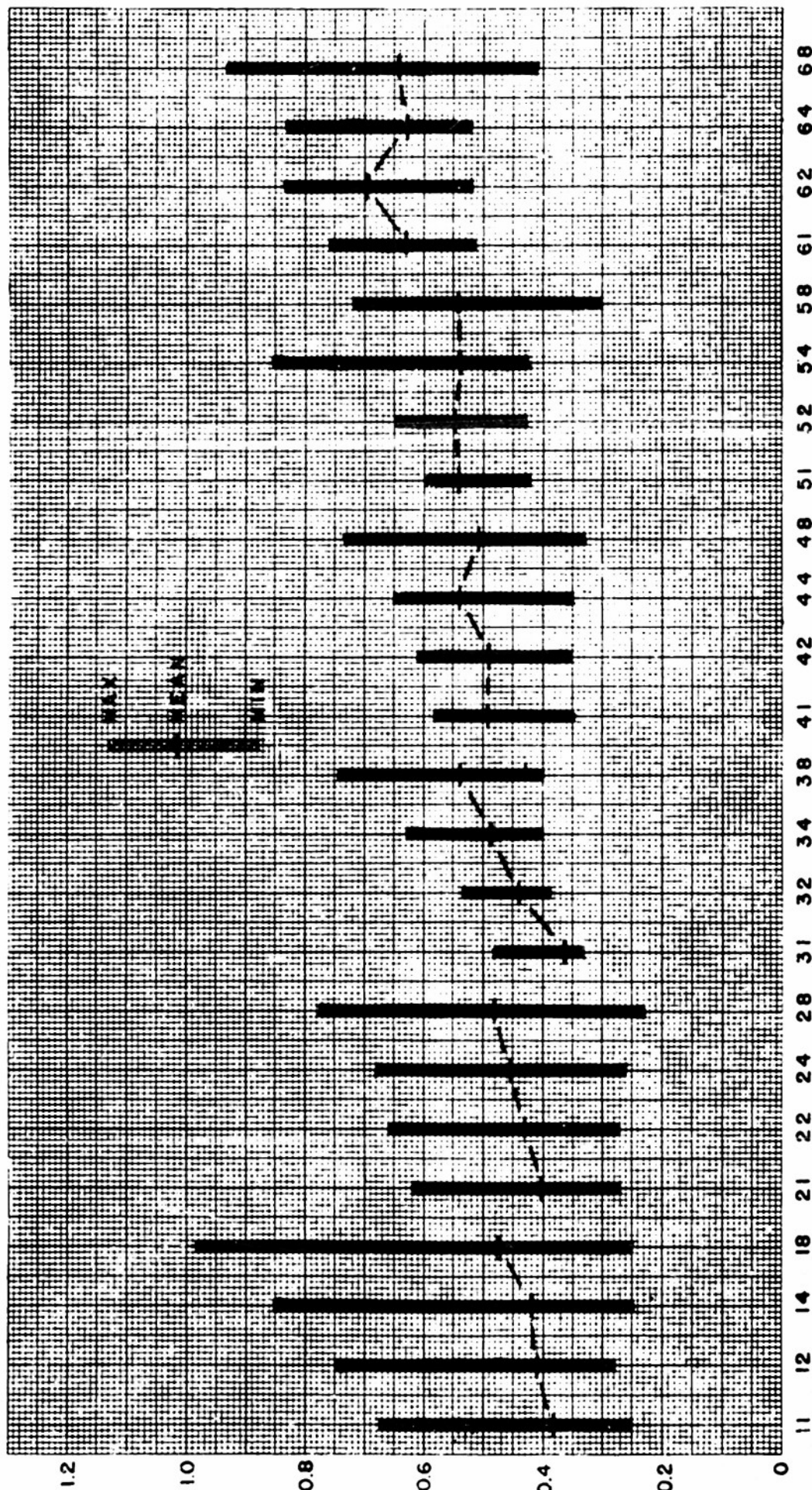
HORIZONTAL SCALE IDENTIFICATION

First Digit	Second Digit	Time Av. (Sec)
Anemometer Height	Space Av.	
1 25 ft.	5 25 & 100 ft.	1
2 50 ft.	6 50 & 150 ft.	2
3 100 ft.		4
4 150 ft.		8

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FIG. 13 EFFECT OF ALTITUDE AND AVERAGE ON SLOPE OF REGRESSION LINE,
LOKI BURNOUT DEVIATION VS. WIND SPEED
(University of Michigan Wind Data)

REGRESSION LINE SLOPE, BURNOUT DEVIATION VS. WIND SPEED m/s/mph



HORIZONTAL SCALE IDENTIFICATION

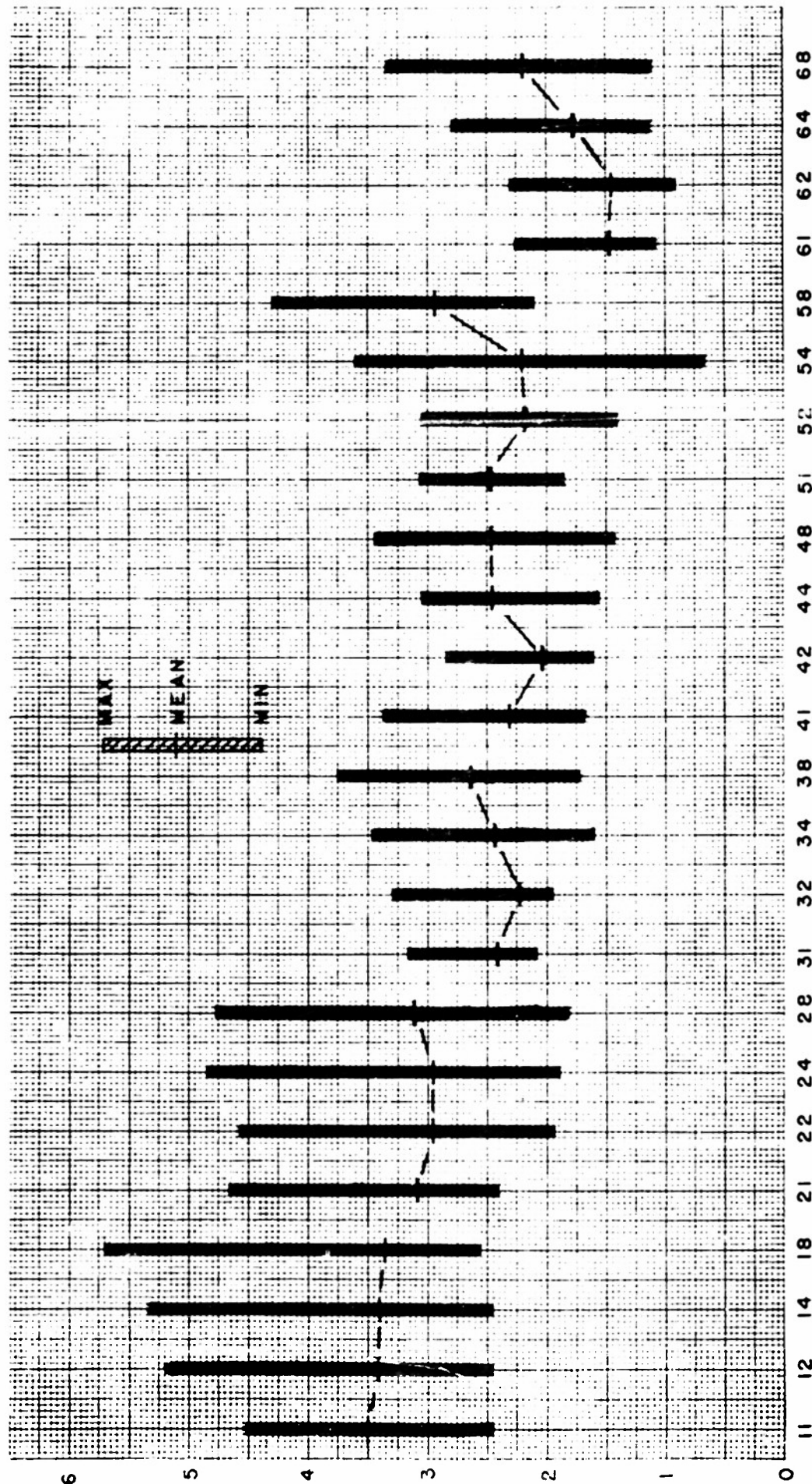
First Digit		Second Digit	
Anemometer Height	Space Av.	Time Av. (Sec)	
1 25 ft.	5 25 & 100 ft.	1	
2 50 ft.	6 50 & 150 ft.	2	
3 100 ft.		4	
4 150 ft.		8	

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FIG. 14 EFFECT OF ALTITUDE AND AVERAGE ON 90 PERCENT CONFIDENCE LIMITS
FOR LOKI BURNOUT DEVIATION
(University of Michigan Wind Data)

90 PERCENT CONFIDENCE LIMITS, BURNOUT DEVIATION VS. WIND SPEED + MILS



HORIZONTAL SCALE IDENTIFICATION

First Digit	Second Digit	Anemometer Height	Space Av.	Time Av. (Sec)
1	25 ft.	5	25 & 100 ft.	1
2	50 ft.	6	50 & 150 ft.	2
3	100 ft.			4
4	150 ft.			8

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The present results are in substantial agreement with those obtained in the preliminary analysis of Ref. 1, and indicate the best choice of altitude and smoothing that would apply to vertical trajectories under the condition that would obtain at the Michigan test site in winter storm conditions.

For an anemometer height of 100 feet and 2-second average, the correction for wind effect in the first 250 feet of a vertical trajectory (about 50% of the total) could be expressed by the relation:

$$-\delta_c = 0.48W + 4$$

where δ = burnout deviation correction, mils

w = 100 foot 2-second average wind velocity, mph,

the 90% confidence limits would be ± 2 mils.

Supplementary data obtained from hot wire surveys up to 1000 feet at White Sands Proving Ground and El Mirage Airport, using equipment described in Ref. 3, is now being analyzed and should permit the establishment of a tentative wind correction correction equation for the entire boost phase.

The effects of horizontal separation between the anemometer and launcher, elevation angle of launching, terrain and weather conditions, must also be considered in the design of a tactical system. These will be investigated in the Signal Corps sponsored project which is now being conducted by North American Instruments, Inc.

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APPENDIX

**STATISTICAL RESULTS OF LOKI BURNOUT DEVIATION
VS. WIND VELOCITY COMPARISON**

The tables which follow present the results of the statistical computations described above, p. 15, and are direct photographs of the original IBM records. The accuracy is no better than 3 figures, the additional figures were allowed to appear as a matter of convenience in machine operation.

The identification code for the various columns as given on p. 15 are repeated below:

p = height of observation code number

m = index for moving average, seconds

k = episode number

N = number of observations

\bar{J} = mean burnout deviation, mils, for episode

\bar{A} = mean wind velocity, mph, for episode

b = regression line slope, mils/mph

s_{y90} = 90% confidence limits, mils

r = correlation coefficient

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p m k	N	$\bar{\delta}$	\bar{A}	b	$S_{y_{30}}$	r	p m k
1 1 01	152	18	401973	38	803947	31	01
1 1 02	227	19	934373	39	026339	31	02
1 1 07	227	24	529629	39	962962	37	07
1 1 08	184	23	998369	37	620108	37	08
1 1 09	200	20	318500	30	773500	30	09
1 1 10	155	22	876129	36	141290	36	10
1 1 11	91	21	609890	33	123626	33	11
1 1 16	142	18	761267	30	643661	30	16
1 1 17	115	22	553043	35	077391	35	17
1 1 18	114	22	219298	33	531578	33	18
1 1 19	43	18	365116	28	067441	28	19
1 1 20	127	22	246456	35	396062	35	20
1 1 21	90	20	978888	32	227777	32	21
1 1 22	93	24	375268	36	915053	36	22
1 1 23	115	22	133043	34	770434	34	23
1 1 24	55	22	290909	35	580000	35	24
1 1 25	120	17	281666	28	977500	28	25
1 1 26	116	15	011206	26	203448	26	26
1 1 27	80	16	117500	26	157500	26	27
1 1 28	119	17	319327	28	700000	28	28
1 1 29	129	20	050387	32	242635	32	29
1 1 30	77	23	925974	36	853246	36	30
1 1 31	118	22	050000	33	775423	33	31
1 1 32	117	19	694871	33	456410	33	32
1 1 33	111	17	792792	29	000000	29	33
1 1 34	38	14	471052	24	239473	24	34
1 1 35	53	15	956603	26	596226	26	35
1 1 36	114	16	224561	26	524561	26	36
1 1 37	113	21	899115	32	083185	32	37
1 1 38	61	20	645901	32	540983	32	38

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p m k	N	$\bar{\delta}$	\bar{A}	b	Sy ₉₀	r	ρ	m k
1 1	151	18	438410	0	5182796	8084899	1	01
1 1	223	19	934977	0	2669432	5380295	1	02
1 1	26	24	596153	0	351375	7239712	1	07
1 1	183	24	021311	0	511611	5338422	1	08
1 1	199	20	296984	0	4016338	6434189	1	09
1 1	154	22	812727	0	4293502	7528113	1	10
1 1	50	22	656666	0	2917832	4710685	1	11
1 1	141	18	763829	0	3270586	6285578	1	16
1 1	113	22	523684	0	2831245	4850075	1	17
1 1	113	22	402380	0	2854644	5407539	1	17
1 1	42	18	273015	0	0784893	1232166	1	19
1 1	126	22	998876	0	3216400	7490947	1	20
1 1	89	20	347826	0	4415197	6571905	1	21
1 1	114	22	107894	0	4012063	7331157	1	22
1 1	54	22	429629	0	2706983	6697481	1	23
1 1	119	17	292436	0	4467008	5503548	1	24
1 1	115	14	992173	0	4373757	7902319	1	25
1 1	7	16	196202	0	4373757	6757290	1	26
1 1	118	17	325423	0	3624354	5578091	1	27
1 1	128	17	061716	0	5377464	7980858	1	28
1 1	76	23	965789	0	5377464	7112711	1	29
1 1	117	22	033333	0	4656183	7112711	1	30
1 1	116	19	683620	0	6018220	67153307	1	31
1 1	110	17	786363	0	3656028	75537307	1	32
1 1	37	14	464864	0	3378415	6479284	1	33
1 1	52	24	254054	0	8204932	6600850	1	34
1 1	113	15	961538	0	1341477	7169922	1	35
1 1	112	16	285840	0	7523865	1787909	1	36
1 1	60	21	914285	0	4316750	8100060	1	37
1 1	60	20	561666	0	3962969	6088897	1	38
1 1	60	20	455000	0	3962969	5670305	1	38

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p m k	N	$\bar{\delta}$	\bar{A}	b	Sy ₅₀	r	p m k
1 4 0 1	149	18	28	0	3	7986069	1 4 0 1
1 4 0 2	221	19	31	5378935	3	5201446	1 4 0 2
1 4 0 7	24	24	40	2797660	1	8562992	1 4 0 7
1 4 0 8	181	24	37	8500091	4	784851	1 4 0 8
1 4 0 9	197	20	30	3127410	4	5078233	1 4 0 9
1 4 1 0	152	22	36	707734	3	6238836	1 4 1 0
1 4 1 1	88	21	36	757360	3	714187	1 4 1 1
1 4 1 6	139	18	30	163157	3	7387583	1 4 1 6
1 4 1 7	112	22	33	275000	3	4960701	1 4 1 7
1 4 1 8	111	22	35	3418285	3	5987934	1 4 1 8
1 4 1 9	40	18	33	3329951	3	5205151	1 4 1 9
1 4 2 0	124	22	35	3452217	3	5242754	1 4 2 0
1 4 2 1	87	21	33	2959241	3	803744	1 4 2 1
1 4 2 2	90	24	35	0606041	3	508688	1 4 2 2
1 4 2 3	112	22	35	3141569	3	786480	1 4 2 3
1 4 2 4	117	22	32	5009927	3	7084947	1 4 2 4
1 4 2 5	113	17	36	5283344	2	6661445	1 4 2 5
1 4 2 6	116	17	34	4612355	2	7710487	1 4 2 6
1 4 2 7	117	16	35	2523129	4	6932569	1 4 2 7
1 4 2 8	116	17	35	4770931	2	4630056	1 4 2 8
1 4 2 9	126	20	26	4555805	3	7816312	1 4 2 9
1 4 3 0	74	24	26	3567564	3	5981736	1 4 3 0
1 4 3 1	115	21	37	572912	3	7963009	1 4 3 1
1 4 3 2	114	19	32	5729306	3	910603	1 4 3 2
1 4 3 3	108	17	32	3358190	3	6969398	1 4 3 3
1 4 3 4	35	14	33	6564119	3	6122177	1 4 3 4
1 4 3 5	50	16	33	3876673	3	7299101	1 4 3 5
1 4 3 6	111	16	32	3758067	3	62774202	1 4 3 6
1 4 3 7	110	21	29	0540505	3	7210924	1 4 3 7
1 4 3 8	58	20	32	1609099	3	1947797	1 4 3 8
				8499064	3	8730651	
				4842654	3	6161995	
				3778154	3	4918028	

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p m k	N	$\bar{\delta}$	\bar{A}	b	$S_{4\infty}$	r	p	m	k
1 1 1	145	18	565895	39	008965	7529897	1	8	001
1 1 1	217	19	543317	30	967741	4627271	1	8	002
1 1 1	177	24	068561	37	819774	0000000	1	8	008
1 1 1	193	20	142487	30	523160	0000000	1	8	009
1 1 1	140	22	906756	36	198648	4151520	1	8	008
1 1 1	184	21	735714	33	263095	5524350	1	8	010
1 1 1	135	18	735555	30	583703	7141287	1	8	011
1 1 1	108	22	391666	35	084259	5005269	1	8	016
1 1 1	107	22	410280	33	632710	5313945	1	8	017
1 1 1	136	19	694444	28	072777	4200011	1	8	018
1 1 1	120	22	459166	35	885833	3201745	1	8	019
1 1 1	83	21	106024	32	204819	6141147	1	8	020
1 1 1	86	24	343023	36	673255	6625357	1	8	021
1 1 1	108	21	931481	34	863388	6784269	1	8	022
1 1 1	48	22	477083	36	337300	7147729	1	8	023
1 1 1	113	17	385840	29	032743	3056662	1	8	024
1 1 1	109	14	644036	26	142201	7523553	1	8	025
1 1 1	73	16	697260	25	968493	6366686	1	8	026
1 1 1	112	17	164285	28	764285	6805529	1	8	027
1 1 1	122	20	191803	32	101539	7757029	1	8	028
1 1 1	70	23	927142	37	427142	6398907	1	8	029
1 1 1	111	21	691891	33	752792	7940939	1	8	030
1 1 1	110	19	746153	32	509090	6205204	1	8	031
1 1 1	104	17	746153	29	244230	6205204	1	8	032
1 1 1	31	14	161290	24	177419	6897394	1	8	033
1 1 1	46	16	067391	26	439130	6227726	1	8	034
1 1 1	107	16	671952	26	357943	2554539	1	8	035
1 1 1	106	21	955943	32	165094	9278645	1	8	036
1 1 1	54	20	244444	32	140740	6571454	1	8	037
1 1 1						3044165	1	8	038

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p m k		N		δ		A		b		S ₄₃₀		r		p m k	
1	01	152	18	401973	30	997368	00	4635156	3	1862037	3	8509333	1	01	
1	02	124	19	934373	33	412946	00	3059812	3	1954937	3	6018881	1	02	
1	03	245	24	051020	39	126122	00	3898441	3	746775	3	6227334	1	03	
1	04	161	23	554037	38	726096	00	4053040	3	555044	3	8402051	1	04	
1	05	231	24	049783	39	477922	00	4158410	3	118592	3	7870042	1	05	
1	06	58	20	656896	34	513793	00	5140887	3	211185	3	7597477	1	06	
1	07	195	25	320512	41	769743	00	3535244	3	019136	3	7028496	1	07	
1	08	184	23	998369	40	243478	00	3660894	3	649816	3	7311327	1	08	
1	09	200	20	318500	32	824000	00	4083720	3	634087	3	7201528	1	09	
1	10	155	22	876129	38	451612	00	4202415	3	160501	3	8165507	1	10	
1	11	91	21	761267	36	024175	00	3537862	3	086002	3	5682772	1	11	
1	12	142	18	553043	31	261971	00	3485456	3	726429	3	7374043	1	12	
1	13	115	22	319298	37	320869	00	3823910	3	729150	3	6109853	1	13	
1	14	114	22	365116	37	477192	00	3158140	3	323900	3	6714401	1	14	
1	15	43	18	246456	30	316279	00	3602523	3	638798	3	6557905	1	15	
1	16	127	22	978888	37	648031	00	3453788	3	328210	3	8150610	1	16	
1	17	90	20	375268	34	215555	00	4338653	3	457706	3	8456795	1	17	
1	18	93	24	133043	39	50537	00	4709212	3	424973	3	8278560	1	18	
1	19	115	22	290909	35	834782	00	4024683	3	028449	3	6652497	1	19	
1	20	55	22	281666	37	580000	00	3163873	3	462785	3	6410649	1	20	
1	21	120	15	011206	28	603333	00	4121933	3	614387	3	7571013	1	21	
1	22	116	16	117500	25	121551	00	4000587	3	069553	3	6813314	1	22	
1	23	80	17	319327	28	383750	00	3426612	4	093162	3	6263627	1	23	
1	24	119	17	925974	36	831932	00	3268018	3	009227	3	8829903	1	24	
1	25	129	20	050387	31	644186	00	5069095	3	664754	3	7798476	1	25	
1	26	77	23	925974	36	906493	00	3901990	3	964754	3	7503654	1	26	
1	27	117	19	694871	32	136440	00	4912793	3	065236	3	7871171	1	27	
1	28	111	17	792792	32	053846	00	3876665	3	617912	3	7306377	1	28	
1	29	38	14	471052	23	210810	00	2702952	3	649776	3	6702255	1	29	
1	30	53	15	955603	25	952663	00	4068766	3	248419	3	7708208	1	30	
1	31	114	16	224561	25	926415	00	3209426	3	422155	3	7420826	1	31	
1	32	113	21	899115	32	151754	00	6204667	3	975192	3	8304852	1	32	
1	33	61	20	645901	32	983185	00	3957636	3	461746	3	6779084	1	33	
1	34					937704	00	3797544	2	996935	2	7170005	1	34	

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p m k		N		$\bar{\delta}$	\bar{A}	b		$S_{y_{30}}$	r		p m k	
3	2	3	2	3	2	3	2	3	2	3	2	3
01	01	151	151	19	31	438419	043046	3	8523853	01	01	01
03	03	244	244	24	39	085655	155327	3	6186603	03	03	03
05	05	160	160	23	38	546875	779375	3	8453205	05	05	05
07	07	230	230	24	39	053913	490434	3	7900917	07	07	07
09	09	57	57	20	34	619298	543859	3	8181744	09	09	09
11	11	194	194	25	34	331958	803092	3	7114289	11	11	11
13	13	183	183	24	40	021311	284699	3	7321613	13	13	13
15	15	199	199	20	32	296584	759296	3	7161769	15	15	15
17	17	154	154	22	38	872727	479870	3	8204331	17	17	17
19	19	90	90	21	36	556666	067777	3	6151571	19	19	19
21	21	141	141	18	31	763829	292198	3	7541994	21	21	21
23	23	114	114	22	37	523684	326315	3	6350129	23	23	23
25	25	113	113	22	37	243362	539823	3	6784384	25	25	25
27	27	42	42	22	30	402380	321428	3	6406269	27	27	27
29	29	126	126	20	34	273015	723809	3	8178564	29	29	29
31	31	89	89	24	39	998876	243320	3	8278991	31	31	31
33	33	92	92	24	39	347826	532608	3	8677088	33	33	33
35	35	114	114	22	37	107894	846491	3	7278310	35	35	35
37	37	54	54	22	35	296229	712962	3	6013294	37	37	37
39	39	119	119	17	28	292436	626890	3	7774204	39	39	39
41	41	115	115	14	26	992173	140000	3	6989675	41	41	41
43	43	79	79	16	25	196202	92405	3	6202251	43	43	43
45	45	118	118	17	28	325423	855932	3	8981598	45	45	45
47	47	76	76	20	31	061718	667187	3	7841149	47	47	47
49	49	117	117	23	37	965799	053947	3	971697	49	49	49
51	51	116	116	22	33	033333	147863	3	8068498	51	51	51
53	53	110	110	19	32	583620	050000	3	7340230	53	53	53
55	55	137	137	17	29	786363	277272	3	936951	55	55	55
57	57	52	52	15	26	464864	870270	3	6751752	57	57	57
59	59	113	113	16	25	561538	013461	3	7654976	59	59	59
61	61	112	112	21	32	285840	152212	3	3059556	61	61	61
63	63	60	60	20	32	914285	994642	3	8640595	63	63	63
65	65					516166	851666	3	6982213	65	65	65
67	67							3	7270463	67	67	67

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p m k	N	$\bar{\delta}$	\bar{A}	b	$S_{y_{50}}$	r	p m k
1	149	18	31	0	287800	8357025	1
0 0 2	1221	19	33	0	934720	6118490	0 0 2
0 0 3	2242	24	39	0	732718	5857275	0 0 3
0 0 4	158	23	38	0	621585	8317664	0 0 4
0 0 5	228	24	39	0	248318	7685737	0 0 5
0 0 6	155	20	34	0	881979	8281503	0 0 6
0 0 7	192	25	41	0	157141	6658968	0 0 7
0 0 8	181	24	40	0	849068	6921376	0 0 8
0 0 9	157	20	32	0	793767	6794745	0 0 9
0 1 0	152	22	38	0	354008	7934672	0 1 0
0 1 1	188	21	36	0	690261	6782240	0 1 1
0 1 6	139	18	31	0	679419	7522624	0 1 6
0 1 7	112	22	37	0	619546	6303194	0 1 7
0 1 8	111	22	37	0	436673	6388006	0 1 8
0 2 0	140	18	30	0	069489	4898713	0 2 0
0 2 1	124	22	37	0	399803	7941021	0 2 1
0 2 2	87	21	34	0	041995	8966977	0 2 2
0 2 3	90	24	39	0	957038	8908178	0 2 3
0 2 4	112	22	35	0	769871	7408725	0 2 4
0 2 5	117	22	37	0	230182	4873431	0 2 5
0 2 6	113	17	28	0	454956	7927807	0 2 6
0 2 7	77	16	25	0	971088	6983261	0 2 7
0 2 8	116	17	28	0	756496	6462735	0 2 8
0 3 0	126	20	31	0	747891	9051700	0 3 0
0 3 1	74	24	37	0	7622002	7046083	0 3 1
0 3 2	115	21	33	0	181781	8091553	0 3 2
0 3 3	114	19	32	0	899588	7203262	0 3 3
0 3 4	108	17	29	0	682145	6724573	0 3 4
0 3 5	35	14	23	0	662010	7042145	0 3 5
0 3 6	50	16	26	0	539234	7992261	0 3 6
0 3 7	110	16	25	0	612586	8915683	0 3 7
0 3 8	158	20	32	0	139279	7018167	0 3 8
					350251	6643957	
					982927		

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p m k	N	$\bar{\delta}$	\bar{A}	b	54 ₃₀	r	p m k
01	145	18	31	0	3	78533056	01
02	217	19	33	0	3	5688950	02
03	238	24	39	0	3	5312583	03
04	154	23	38	0	3	7967023	04
05	224	24	39	0	3	7112844	05
06	51	20	34	0	3	8029291	06
07	188	25	41	0	3	6173571	07
08	177	24	40	0	3	6160856	08
09	193	20	32	0	3	6167949	09
10	148	22	38	0	3	7367053	10
11	84	21	36	0	3	6291970	11
16	135	18	31	0	3	7003024	16
17	108	22	37	0	3	6579240	17
18	107	22	37	0	3	5024287	18
19	36	18	30	0	3	1996292	19
20	120	22	38	0	3	7312551	20
21	83	21	34	0	3	8666546	21
22	86	24	39	0	3	7970538	22
23	108	21	35	0	3	7425787	23
24	48	22	38	0	3	2713005	24
25	113	17	28	0	3	7719676	25
26	103	14	26	0	3	6811951	26
27	73	15	25	0	3	7018547	27
28	112	17	28	0	3	8921524	28
29	122	20	31	0	3	7095538	29
30	70	23	37	0	3	7583581	30
31	111	21	32	0	3	8107163	31
32	110	19	32	0	3	7118562	32
33	104	17	29	0	3	6581064	33
34	31	14	23	0	3	6053602	34
35	46	16	26	0	3	2017818	35
36	107	16	25	0	3	8989802	36
37	106	21	32	0	3	6927510	37
38	54	20	32	0	3	3920305	38

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p m k	N	$\bar{\delta}$	\bar{A}	b	S ₄₀₀	r	p m k		
1 0 2	152	401973	903947	0	344606	9077550	3 3	0 1	
1 0 3	245	934373	126785	3467764	3477168	7428504	3 3	1 0 3	
1 0 4	245	051020	40167080	0	019750	7760995	3 3	1 0 4	
1 0 5	161	554037	40167080	0	353942	8662633	3 3	1 0 5	
1 0 6	231	049783	401754112	0	408266	8792154	3 3	1 0 6	
1 0 7	58	656896	34179310	0	236219	7533852	3 3	1 0 7	
1 0 8	195	320512	42693333	0	020056	7773312	3 3	1 0 8	
1 0 9	184	998369	39780978	0	236143	7962938	3 3	1 0 9	
1 1 0	200	318500	32823500	0	339789	8004161	3 3	1 1 0	
1 1 1	155	876129	38009677	0	356145	8843169	3 3	1 1 1	
1 1 2	91	609890	34945054	0	322301	7400657	3 3	1 1 2	
1 1 3	142	751267	31676760	0	350242	7870118	3 3	1 1 3	
1 1 4	115	553043	37055652	0	309747	7423736	3 3	1 1 4	
1 1 5	114	219298	37669298	0	701214	7983212	3 3	1 1 5	
1 1 6	43	365116	29616279	0	27404	6082588	3 3	1 1 6	
1 1 7	127	246456	38232283	0	245353	8291632	3 3	1 1 7	
1 1 8	90	978888	34563333	0	305784	8778227	3 3	1 1 8	
1 1 9	93	375268	41065591	0	222426	8577376	3 3	1 1 9	
1 2 0	115	133043	36474732	0	252645	7495213	3 3	1 2 0	
1 2 1	55	290909	36865454	0	781240	8334484	3 3	1 2 1	
1 2 2							3 3	1 2 2	
1 2 3							3 3	1 2 3	
1 2 4							3 3	1 2 4	

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p m k	N	$\bar{\delta}$	\bar{A}	b	S ₄₀₀	r	p m k
01	151	18	438410	0	2	9162588	3
02	223	19	934977	0	2	7828251	3
03	244	24	085655	0	2	7994412	3
04	160	23	546875	0	1	8104174	3
05	230	24	053913	0	2	8556325	3
06	57	20	619298	0	2	7838092	3
07	194	25	331958	0	3	7903575	3
08	183	24	021311	0	3	8219333	3
09	193	20	296984	0	3	7712587	3
10	154	22	872727	0	2	8982339	3
11	90	21	656666	0	2	7829549	3
16	141	18	763829	0	2	8253998	3
17	114	22	523684	0	2	7849553	3
18	113	22	243362	0	2	8253623	3
19	42	18	402380	0	2	6845891	3
20	126	22	273015	0	2	8525102	3
21	89	20	998876	0	2	8924078	3
22	92	24	347826	0	2	8673738	3
23	114	22	107894	0	2	7595680	3
24	54	22	429629	0	2	8420723	3

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p m k	N	$\bar{\delta}$	\bar{A}	b	S ₄₅₀	r	p m k
01	149	510067	30 964429	0 4377103	2 596173	9010788	01
02	221	934841	33 102262	0 4429923	2 247738	7955916	02
03	242	143801	40 507024	0 5204809	2 831230	7884461	03
04	158	522784	40 172784	0 5894528	1 640436	9375215	04
05	229	060087	40 768421	0 5014784	2 365399	8847916	05
06	55	560000	34 292727	0 4223245	2 096735	7810516	06
07	192	362500	42 796875	0 4798105	2 770134	8219256	07
08	181	050826	39 875690	0 4462753	3 037315	7559819	08
09	197	244162	32 717258	0 4123475	3 506519	7349232	09
10	152	869078	37 990789	0 4596477	2 505906	8906301	10
11	88	718181	35 126136	0 4864582	2 315696	7745266	11
16	139	758273	31 694964	0 6393724	2 086912	8582842	16
17	112	467857	37 074107	0 5123763	1 919800	8223672	17
18	111	290930	37 685585	0 4623746	2 625315	8090228	18
19	40	460000	29 947500	0 5733296	2 227411	7744558	19
20	124	331451	39 402419	0 5161982	1 590335	8636565	20
21	87	051724	34 577011	0 4694172	2 049132	8959324	21
22	90	294444	40 991111	0 4410437	2 484916	8168029	22
23	112	062500	36 466964	0 4259549	2 692659	7574241	23
24	52	467307	36 705769	0 5426938	2 820658	8130094	24

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p m k	N	δ	\bar{A}	b	Sy ₉₀	r	p m k
01	145	665896	31 064137	0	268320	8292833	01
02	117	943317	33 047684	0	486710	7439137	02
03	238	249579	40 554621	0	025811	7300026	03
04	154	479870	40 198701	0	025105	9034794	04
05	224	056250	40 733928	0	839741	8312954	05
06	51	317647	34 266666	0	254410	6745461	06
07	188	354255	42 858510	0	958018	7166124	07
08	177	068361	39 950847	0	325114	7861047	08
09	220	142487	32 524352	0	4789867	6650118	09
10	148	906756	32 982837	0	4276919	8664302	10
11	84	725714	35 276190	0	4809633	7365954	11
12	135	735555	31 720000	0	5507070	9043703	12
13	108	391666	37 064814	0	7592762	8390973	13
14	107	410280	37 682242	0	5996420	7240453	14
15	36	694444	30 250000	0	4454284	6984213	15
16	120	459166	38 580000	0	7134125	6315285	16
17	83	106024	34 607228	0	5256860	8560029	17
18	86	243023	40 933720	0	4794133	6216327	18
19	108	931481	36 431481	0	4000576	7421054	19
20	108	931481	36 431481	0	4744592	7421054	20
21	48	477083	36 729166	0	6465799	7064143	21
22							22
23							23
24							24

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p m k		N		$\bar{\delta}$		\bar{A}		b		54 ₃₀		r		p m k	
4	1	01	152	18	401973	31	490131	0	4901420	2	206660	9314820	4	1	01
4	1	02	224	19	934375	34	063839	0	4526163	2	2533706	7931191	4	1	02
4	1	03	161	23	554037	40	262111	0	0000000	2	0000000	0000000	4	1	03
4	1	04	231	24	049783	41	167532	0	0820339	2	503823	8471225	4	1	04
4	1	05	58	20	656896	35	362068	0	5054607	2	442980	8754592	4	1	05
4	1	06	195	25	320512	43	434358	0	4888948	2	697963	8664287	4	1	06
4	1	07	184	23	998369	41	303260	0	4455650	2	525188	8037558	4	1	07
4	1	08	200	20	318500	35	309000	0	4212529	2	720802	8610143	4	1	08
4	1	09	155	22	876129	39	373528	0	4251178	2	669370	8603919	4	1	09
4	1	10	191	21	609890	36	919780	0	5061133	2	447308	8945304	4	1	10
4	1	11	142	18	761267	31	444366	0	3833153	2	359799	7772386	4	1	11
4	1	12	115	22	253043	37	469565	0	4309898	2	432803	7968368	4	1	12
4	1	13	114	22	219298	37	617548	0	4403473	2	996696	8151399	4	1	13
4	1	14	120	18	365116	32	495348	0	4711264	2	390585	8461294	4	1	14
4	1	15	120	22	245456	38	674013	0	3661346	2	102088	7989528	4	1	15
4	1	16	93	20	978888	36	522222	0	3583721	2	046153	8606570	4	1	16
4	1	17	115	22	375258	42	569892	0	5202866	2	164689	8903067	4	1	17
4	1	18	120	22	133043	37	895652	0	5610894	2	887425	8996577	4	1	18
4	1	19	120	22	290909	38	416363	0	3945969	2	313433	8309450	4	1	19
4	1	20	116	17	281666	39	493333	0	3745160	2	396765	7379142	4	1	20
4	1	21	80	15	011206	35	849137	0	4633137	2	186508	8375395	4	1	21
4	1	22	119	16	117500	27	748750	0	4945472	2	855422	8962934	4	1	22
4	1	23	129	17	319327	27	026050	0	3573365	2	797291	8462881	4	1	23
4	1	24	129	20	050387	34	880620	0	4782916	2	307019	9330071	4	1	24
4	1	25	77	23	925974	42	305194	0	4949917	2	423706	9456369	4	1	25
4	1	26	118	23	050000	42	305194	0	5297366	2	158878	8765405	4	1	26
4	1	27	111	17	694871	38	305932	0	4080137	2	255393	8910820	4	1	27
4	1	28	53	17	792792	33	263963	0	4080137	2	101695	8363911	4	1	28
4	1	29	111	14	471052	33	447368	0	3825140	2	278508	7698942	4	1	29
4	1	30	113	15	956603	29	066037	0	4165706	2	040349	8159667	4	1	30
4	1	31	113	16	224561	29	391238	0	4165706	2	589497	8251208	4	1	31
4	1	32	113	21	899115	38	435398	0	5429738	2	169773	9318482	4	1	32
4	1	33	161	20	645901	35	460655	0	4012159	2	895292	8875148	4	1	33
4	1	34							4495469	1	915196	8952992	4	1	34

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p m k	N	$\bar{\delta}$	\bar{A}	b	Sy ₉₀	r	p m k
44	81	18	338499	31	598436	0	81
44	03	23	546875	40	278325	0	03
44	04	24	053913	41	189130	0	04
44	05	20	619298	35	342105	0	05
44	06	25	331958	43	493814	0	06
44	07	24	021311	41	389071	0	07
44	08	20	296984	35	309045	0	08
44	09	22	872727	38	401948	0	09
44	10	21	636666	36	991111	0	10
44	11	18	763829	31	469503	0	11
44	16	22	523684	38	468421	0	16
44	17	22	243362	37	631858	0	17
44	18	18	402380	32	626190	0	18
44	19	22	273015	38	731746	0	19
44	20	20	998876	36	525842	0	20
44	21	24	147826	42	550000	0	21
44	22	22	107894	37	893659	0	22
44	23	22	429629	38	283323	0	23
44	24	17	292436	29	547056	0	24
44	25	14	992173	25	904347	0	25
44	26	16	196202	27	841772	0	26
44	27	17	325423	31	050847	0	27
44	28	20	061718	34	866406	0	28
44	29	23	965789	42	457294	0	29
44	30	22	033333	38	325641	0	30
44	31	19	683620	33	326724	0	31
44	32	17	786363	32	272727	0	32
44	33	14	464864	25	313513	0	33
44	34	15	961538	29	190384	0	34
44	35	16	285840	29	422433	0	35
44	36	21	914285	38	462500	0	36
44	37	20	561666	35	343333	0	37
44	38						38

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p m k		N	δ	\bar{A}	b	S_{40}	r	p m k	
4	4	149	18	51 0067	0	1	9466297	4	01
4	4	150	19	53 4841	0	2	8140913	4	02
4	4	151	20	52 2784	0	3	0000000	4	03
4	4	152	21	56 0087	0	4	8331170	4	04
4	4	153	22	56 0000	0	5	8658432	4	05
4	4	154	23	56 2500	0	6	8723994	4	06
4	4	155	24	05 0828	0	7	7674075	4	07
4	4	156	25	05 0828	0	8	8507120	4	08
4	4	157	26	24 4162	0	9	8106189	4	09
4	4	158	27	86 9078	0	10	9068235	4	10
4	4	159	28	71 8181	0	11	7214480	4	11
4	4	160	29	75 8273	0	12	8082039	4	12
4	4	161	30	46 7857	0	13	7742910	4	13
4	4	162	31	29 0990	0	14	8377516	4	14
4	4	163	32	46 0000	0	15	8522938	4	15
4	4	164	33	33 1451	0	16	8684946	4	16
4	4	165	34	05 1724	0	17	9907717	4	17
4	4	166	35	29 4444	0	18	7865157	4	18
4	4	167	36	06 2500	0	19	6260135	4	19
4	4	168	37	46 7307	0	20	8963460	4	20
4	4	169	38	30 6837	0	21	8560549	4	21
4	4	170	39	93 3628	0	22	9091232	4	22
4	4	171	40	36 1038	0	23	6009137	4	23
4	4	172	41	29 3965	0	24	5473353	4	24
4	4	173	42	10 6349	0	25	9461140	4	25
4	4	174	43	01 2162	0	26	5326473	4	26
4	4	175	44	97 4334	0	27	8782609	4	27
4	4	176	45	68 3333	0	28	8222837	4	28
4	4	177	46	76 9444	0	29	8294103	4	29
4	4	178	47	40 5714	0	30	8984103	4	30
4	4	179	48	01 4000	0	31	8257173	4	31
4	4	180	49	42 4324	0	32	5426950	4	32
4	4	181	50	96 7272	0	33	8426532	4	33
4	4	182	51	40 3448	0	34	8331480	4	34
4	4	183	52		0	35		4	35
4	4	184	53		0	36		4	36
4	4	185	54		0	37		4	37
4	4	186	55		0	38		4	38

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p m k		N		$\bar{\delta}$		\bar{A}		b		S ₄₀₀		r		p m k	
4	4	4	4	18	865896	31	628328	0	5125821	3	314522	9183587	4	88	01
4	4	4	4	24	249579	41	650840	0	5280562	2	483323	7417434	4	88	03
4	4	4	4	23	479870	40	200000	0	5230626	3	603493	8088295	4	88	05
4	4	4	4	24	056250	41	175446	0	5390662	3	073841	7524278	4	88	06
4	4	4	4	25	317647	35	462745	0	7453759	1	411349	8044406	4	88	07
4	4	4	4	25	354255	43	587765	0	6056757	2	687695	8878603	4	88	08
4	4	4	4	24	068361	41	571751	0	4724422	3	153253	7324004	4	88	09
4	4	4	4	22	142487	35	015025	0	5294420	3	383376	8101422	4	88	10
4	4	4	4	22	506756	39	335810	0	5834482	2	542431	7442345	4	88	11
4	4	4	4	21	735714	37	375000	0	5587498	2	241554	8892565	4	88	12
4	4	4	4	18	735555	31	473333	0	5508126	2	426570	7941912	4	88	13
4	4	4	4	22	391666	38	405555	0	5277679	2	068872	8063908	4	88	14
4	4	4	4	22	410280	37	626261	0	5459585	2	611337	7636105	4	88	15
4	4	4	4	16	694444	33	461111	0	4338242	2	538905	8023487	4	88	16
4	4	4	4	22	459166	38	814166	0	7242035	1	662368	8672429	4	88	17
4	4	4	4	21	106024	36	533734	0	6168270	3	370080	8992536	4	88	18
4	4	4	4	24	243023	42	370930	0	5397005	3	460420	8628452	4	88	19
4	4	4	4	21	931481	37	778703	0	5393789	2	370080	6010734	4	88	20
4	4	4	4	22	477082	37	400000	0	5872010	1	851151	7476636	4	88	21
4	4	4	4	14	844036	25	964220	0	5486860	1	253265	9246777	4	88	22
4	4	4	4	16	697260	28	000000	0	6035310	3	858567	8257894	4	88	23
4	4	4	4	17	164285	30	945535	0	3263387	2	247405	8932931	4	88	24
4	4	4	4	20	191803	34	781147	0	5147064	2	041962	9788339	4	88	25
4	4	4	4	23	927142	42	668571	0	5280361	2	831652	9788339	4	88	26
4	4	4	4	21	891891	38	144144	0	7428823	1	489185	9246777	4	88	27
4	4	4	4	19	721818	33	358181	0	4850815	2	817093	8227536	4	88	28
4	4	4	4	17	746153	32	118269	0	5394175	2	176511	8290953	4	88	29
4	4	4	4	14	161290	24	751612	0	6143493	1	079502	8190265	4	88	30
4	4	4	4	16	067391	29	710869	0	5778973	2	604095	8646054	4	88	31
4	4	4	4	16	671962	29	456074	0	6510997	2	514640	7151036	4	88	32
4	4	4	4	21	959433	38	360377	0	4442802	2	856788	9222002	4	88	33
4	4	4	4	20	244444	34	622222	0	5127647	2	690517	7906190	4	88	34
4	4	4	4	20	244444	34	622222	0	5127647	2	690517	7349885	4	88	35

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p m k	N	$\bar{\delta}$	\bar{A}	b	$S_{y_{90}}$	r	p m k
5 1 01	152	18	401973	0	5407305	5345421	5 1 01
5 1 02	227	19	934375	0	4316620	7651030	5 1 02
5 1 03	184	24	529629	0	6799641	8372847	5 1 03
5 1 04	200	22	998369	0	4982545	8100541	5 1 04
5 1 05	155	20	318500	0	5564413	8807824	5 1 05
5 1 06	91	22	876129	0	5138717	9106282	5 1 06
5 1 07	142	21	609890	0	5901999	7864651	5 1 07
5 1 08	115	22	761267	0	5528859	8353400	5 1 08
5 1 09	114	22	553043	0	5447951	7825519	5 1 09
5 1 10	43	22	212296	0	4816817	7995373	5 1 10
5 1 11	127	18	365116	0	5919374	6711797	5 1 11
5 1 12	90	22	246456	0	4750346	8918245	5 1 12
5 1 13	93	20	978888	0	5709016	8904727	5 1 13
5 1 14	115	24	375268	0	5295127	8858713	5 1 14
5 1 15	55	22	133043	0	5423239	8493659	5 1 15
5 1 16	55	22	290909	0	4612386	7917535	5 1 16

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p m k	N	$\bar{\delta}$	\bar{A}	b	Sy ₉₀	r	p m k
555	155	18	32	0	3	9305649	555
555	156	24	40	0	1	7803482	555
555	157	24	38	0	3	5272465	555
555	158	20	31	0	3	8152664	555
555	159	22	37	0	2	5408136	555
555	160	21	34	0	2	5070837	555
555	161	18	31	0	2	6245221	555
555	162	22	36	0	1	8449696	555
555	163	22	35	0	2	8186423	555
555	164	18	28	0	2	8028573	555
555	165	22	36	0	2	6666294	555
555	166	20	33	0	1	8890650	555
555	167	22	36	0	1	5023172	555
555	168	24	39	0	1	5070305	555
555	169	22	35	0	2	8600700	555
555	170	22	36	0	3	7552431	555

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p m k		N		$\bar{\delta}$		\bar{A}		b		S ₄₅₀		r		p m k	
5	4	01	149	16	510067	39	944966	0	5362819	2	572712	9029539	4	01	4
5	4	02	221	19	334841	32	062823	0	4842182	2	404032	7616731	4	02	4
5	4	07	24	24	579166	40	175000	1	1984559	0	661804	5817317	4	07	4
5	4	08	181	24	050828	38	812154	0	5349176	3	264468	7907469	4	08	4
5	4	09	152	20	244162	31	765989	0	5493480	3	180680	7864253	4	09	4
5	4	10	88	22	869078	37	103947	0	5231011	2	609988	8907377	4	10	4
5	4	11	139	21	718181	34	217045	0	7397760	2	018465	8342730	4	11	4
5	4	16	112	18	758273	31	195683	0	5994377	2	337143	8163651	4	16	4
5	4	17	112	22	467857	36	081250	0	6989666	1	736969	8573266	4	17	4
5	4	18	111	22	290390	35	629279	0	4975871	2	680649	7642718	4	18	4
5	4	19	40	18	460000	29	042500	0	7448316	2	857040	5844149	4	19	4
5	4	20	124	22	331451	37	028225	0	4819023	2	040311	8561437	4	20	4
5	4	21	87	21	051724	33	443678	0	6060476	2	112531	8885853	4	21	4
5	4	22	90	24	294444	38	941111	0	5873921	2	042311	8604396	4	22	4
5	4	23	112	22	062500	35	668750	0	5998661	2	194093	8467259	4	23	4
5	4	24	52	22	467307	36	369239	0	4214231	3	618450	6642933	4	24	4

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p m k	N	$\bar{\delta}$	\bar{A}	b	Sy ₃₀	r	p m k
5 8 01	145	18 555896	39 058520	0 4986574	3 295134	8261803	5 8 01
5 8 02	177	24 068361	38 909039	0 4814132	2 237559	6373291	5 8 02
5 8 08	193	20 142487	31 604145	0 0000000	0 0000000	0000000	5 8 08
5 8 09	148	22 906756	37 120270	0 5270295	3 819935	7041357	5 8 09
5 8 10	84	21 735714	34 291666	0 5164097	3 674838	6883808	5 8 10
5 8 11	135	18 735555	31 175552	0 5178080	3 051292	8358501	5 8 11
5 8 16	108	22 391666	36 191666	0 8775298	2 112124	8198486	5 8 16
5 8 17	107	22 410280	35 682242	0 6260285	2 590654	7254779	5 8 17
5 8 18	107	22 410280	35 682242	0 3024166	3 110733	3569896	5 8 18
5 8 19	107	22 410280	35 682242	0 4364775	3 405818	6249805	5 8 19
5 8 20	120	22 459166	37 259166	0 5838877	3 227876	3213388	5 8 20
5 8 21	83	21 106024	33 432530	0 4525792	2 432060	7701592	5 8 21
5 8 22	86	24 243023	38 818604	0 6072300	2 571205	8362409	5 8 22
5 8 23	108	21 931481	35 674074	0 5674842	3 045524	7108215	5 8 23
5 8 24	48	22 477083	36 554166	0 6414528	2 329230	8129303	5 8 24
				0 4101824	4 309112	4989177	5 8 24

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p m k		N		$\bar{\delta}$		\bar{A}		b		54 ₉₀		r		p m k	
1	01	152	18	401973	31	353157	0	5598423	1	570261	5659118	6	1	01	01
6	02	245	19	934373	33	757389	0	61298807	1	556473	9072183	6	1	02	02
6	03	161	24	051020	40	303265	0	7142641	1	452000	9502225	6	1	03	03
6	04	231	23	554037	39	519875	0	5662141	1	375730	9564252	6	1	04	04
6	05	195	20	049783	40	342424	0	6115558	1	476432	9563937	6	1	05	05
6	06	184	25	656896	34	958620	0	6679110	1	141923	9419360	6	1	06	06
6	07	200	23	320512	42	623589	0	6467652	2	1496246	9358004	6	1	07	07
6	08	155	20	998369	40	795652	0	5251777	1	112855	9187063	6	1	08	08
6	09	142	22	318500	34	067000	0	5554740	1	823480	93774138	6	1	09	09
6	10	91	21	876129	38	937419	0	5806092	1	553626	9588912	6	1	10	10
6	11	142	21	609890	36	500000	0	7208274	1	196533	9477419	6	1	11	11
6	12	115	18	761267	31	375352	0	5863049	1	343190	9430099	6	1	12	12
6	13	114	22	553043	37	916260	0	7091000	1	183079	9352717	6	1	13	13
6	14	143	22	319230	37	570175	0	5870257	1	643698	9304338	6	1	14	14
6	15	122	18	365116	31	430232	0	5266056	1	662434	8756555	6	1	15	15
6	16	90	22	246456	38	184251	0	5545226	1	250143	9203726	6	1	16	16
6	17	93	20	978868	35	393333	0	5930329	1	216641	9644697	6	1	17	17
6	18	115	22	375263	41	081720	0	6454385	1	105249	9667657	6	1	18	18
6	19	120	22	133043	36	887826	0	5932664	1	526168	9502106	6	1	19	19
6	20	126	22	290909	38	070909	0	5978382	2	114787	9074401	6	1	20	20
6	21	116	17	281666	29	070000	0	6190649	1	272563	9480924	6	1	21	21
6	22	80	15	011206	26	011206	0	6519048	1	270458	9530042	6	1	22	22
6	23	119	16	117500	26	588750	0	5190865	2	341601	9042998	6	1	23	23
6	24	129	17	319327	29	942857	0	5785135	1	369187	9769275	6	1	24	24
6	25	177	20	050387	33	286821	0	6280527	1	761221	9716731	6	1	25	25
6	26	118	22	225974	39	629870	0	6331115	1	292360	9575898	6	1	26	26
6	27	117	19	050000	35	716610	0	5659048	1	506830	9529230	6	1	27	27
6	28	111	17	792792	32	711111	0	5508453	1	481313	9223582	6	1	28	28
6	29	133	14	471052	30	760360	0	4902844	1	649560	8868727	6	1	29	29
6	30	53	15	956603	27	520754	0	4912734	1	257127	8672615	6	1	30	30
6	31	114	16	224561	27	294736	0	8125554	1	309209	9331497	6	1	31	31
6	32	113	21	899115	35	731858	0	6998395	1	623494	9237783	6	1	32	32
6	33	61	20	645901	34	218672	0	5658415	1	602107	9403430	6	1	33	33
6	34							5240027	1	812803	9067590	6	1	34	34

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p m k	N	\bar{S}	\bar{A}	b	Syso	r	p m k
01	151	18	31	0	1	9632572	01
02	153	19	33	0	1	9632572	02
03	244	24	40	0	1	9632572	03
04	160	24	39	0	1	9632572	04
05	230	24	40	0	1	9632572	05
06	157	20	34	0	0	9632572	06
07	194	25	42	0	1	9632572	07
08	183	24	40	0	2	9632572	08
09	199	20	34	0	2	9632572	09
10	154	22	38	0	1	9632572	10
11	90	21	36	0	1	9632572	11
12	141	18	31	0	1	9632572	12
13	114	22	37	0	1	9632572	13
14	113	22	37	0	1	9632572	14
15	42	18	31	0	1	9632572	15
16	126	22	38	0	1	9632572	16
17	89	20	35	0	1	9632572	17
18	92	24	41	0	1	9632572	18
19	114	22	36	0	1	9632572	19
20	154	17	29	0	1	9632572	20
21	115	14	26	0	1	9632572	21
22	79	16	26	0	2	9632572	22
23	118	17	29	0	1	9632572	23
24	128	20	33	0	1	9632572	24
25	76	22	35	0	1	9632572	25
26	117	22	35	0	1	9632572	26
27	116	19	32	0	1	9632572	27
28	110	17	24	0	1	9632572	28
29	37	14	24	0	1	9632572	29
30	52	15	27	0	1	9632572	30
31	112	21	35	0	1	9632572	31
32	160	20	34	0	1	9632572	32
33							33
34							34
35							35
36							36
37							37
38							38

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p m k	N	$\bar{\delta}$	\bar{A}	b	S_{430}	r	p m k						
01	149	18	19	5651770	2	9412808	003						
02	121	19	33	740271	1	5063143	003						
03	242	24	40	371487	1	9157497	003						
04	158	23	39	566455	1	9174551	005						
05	228	24	40	339912	2	9141555	005						
06	55	20	34	990909	1	9401969	007						
07	192	25	42	658437	2	8545897	007						
08	181	24	40	922651	2	8655601	009						
09	1	20	34	017258	2	797673	009						
10	152	22	38	953947	2	9299430	010						
11	88	21	36	630681	1	9342372	011						
16	139	18	31	387769	1	9323494	016						
17	112	22	37	904464	1	9202462	017						
18	111	22	37	608108	3	8751458	018						
19	140	18	31	705000	2	8056880	019						
20	124	22	38	738709	1	9314554	020						
21	87	21	35	405747	1	9478321	021						
22	90	24	40	956666	1	9052571	022						
23	112	22	36	881250	2	8995960	023						
24	52	22	37	911538	1	9324510	024						
25	117	17	29	130769	1	9481937	025						
26	113	14	26	052212	1	8620594	026						
27	77	16	26	683111	2	9760748	027						
28	126	17	29	958620	2	9462362	028						
29	114	20	33	234126	1	9507965	029						
30	74	24	39	990540	1	9259572	030						
31	115	21	35	716521	1	9082513	031						
32	114	19	32	682456	1	8779193	032						
33	108	17	30	791666	1	8359279	033						
34	35	14	24	480000	1	9371002	034						
35	50	16	27	760000	1	9661779	035						
36	111	16	32	722072	2	8812790	036						
37	110	21	35	750909	2	8047057	037						
38	58	20	33	872586	2	8047057	038						

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p m k	N	δ	\bar{A}	b	$S_{y_{30}}$	r	p m k
001	145	18	31	0	745915	8928417	001
002	121	19	33	0	7473085	8928417	002
003	238	24	40	0	7613909	8928417	003
004	154	23	39	0	5692336	8928417	004
005	224	24	39	0	6110711	8928417	005
006	51	20	35	0	8267083	8928417	006
007	188	25	42	0	7942134	8928417	007
008	172	24	41	0	5212741	8928417	008
009	193	22	33	0	5624537	8928417	009
010	148	22	38	0	5826115	8928417	010
011	84	21	36	0	9621202	8928417	011
012	135	18	31	0	6121976	8928417	012
013	108	22	37	0	9500587	8928417	013
014	107	22	37	0	5386481	8928417	014
015	36	18	32	0	4959349	8928417	015
016	120	22	38	0	5863326	8928417	016
017	83	21	35	0	5991109	8928417	017
018	86	24	40	0	6716946	8928417	018
019	108	22	36	0	6030527	8928417	019
020	113	17	29	0	6372154	8928417	020
021	48	22	37	0	6398294	8928417	021
022	109	14	26	0	2408238	8928417	022
023	73	15	26	0	4539230	8928417	023
024	122	20	33	0	5832257	8928417	024
025	70	23	40	0	6344730	8928417	025
026	111	21	35	0	93335760	8928417	026
027	110	19	32	0	5654575	8928417	027
028	104	17	30	0	5906082	8928417	028
029	31	19	32	0	5506912	8928417	029
030	46	14	24	0	5400664	8928417	030
031	106	16	27	0	6995262	8928417	031
032	37	21	35	0	7554979	8928417	032
033	54	20	33	0	5501634	8928417	033
034					4162820	8928417	034
035							035
036							036
037							037
038							038

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2. Smyth, R. K., Instrumentation for LOKI Wind Measurements, Quarterly Report North American Instruments Inc., Altadena, California, January through March 1953.
3. Sherlock R. H., and Stout, M. B., Wind Structure in Winter Storms, Journal of the Aeronautical Sciences, Vol. 5, 1937.
4. Sherlock, R. H., Variation of Wind Velocity and Gusts with Height, Proceedings of the American Society of Civil Engineers, Vol. 78, April 1952.
5. Baker, Donald L., and Phillips, Warren L., LOKI Wind Response Analysis, Quarterly Report North American Instruments, Inc., Altadena, California, July through September, 1953.

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